

Electromagnetic and Gravitational Wave Signatures of Black Hole Mergers

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Galaxies definitely collide and merge

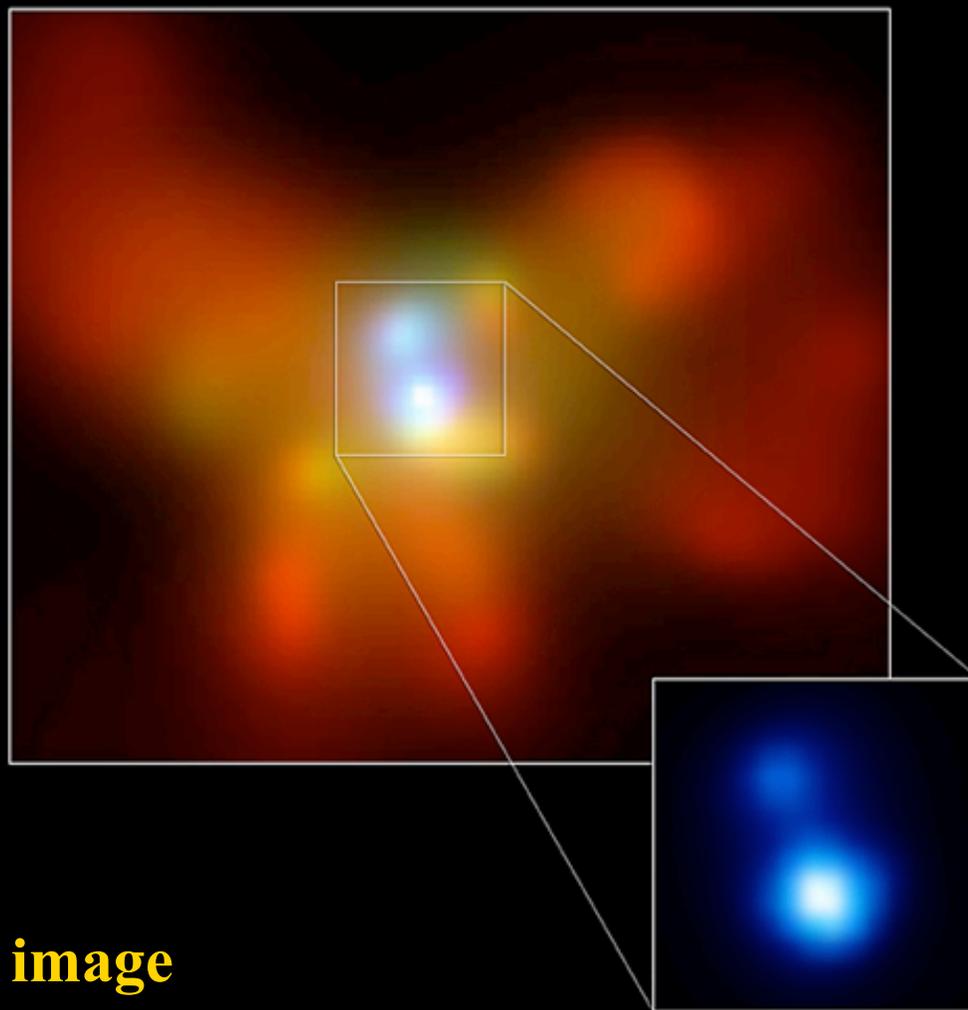


Colliding Galaxies NGC 4038 and NGC 4039

HST • WFPC2

PRC97-34a • ST ScI OPO • October 21, 1997 • B, Whitmore (ST ScI) and NASA

Black hole binary in a galactic nucleus

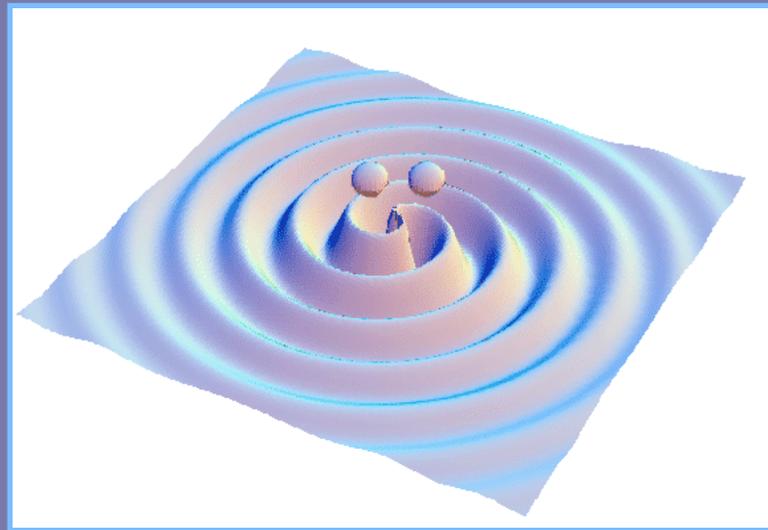


**Chandra X-ray image
of NGC 6240 (Komossa et al. 2003)**

Gravitational Waves by LISA

LISA can detect low-frequency gravitational waves
from super-massive black hole binaries:

sensitive to total mass $(10^4-10^7)/(1+z) M_{\odot}$



Why try hard to find EM counterparts?

- Grav. waves themselves a rich source of info on metric
 - LISA sensitive to BH mass of $\sim(10^4-10^7)/(1+z) M_{\odot}$
- Photons from counterparts: benefits for fundamental physics
 - Hubble diagrams from ‘standard sirens’ ($f(df/dt)^{-1}$; Schutz 1986)
 - $d_L(z)$ from GWs and photons: new test of non-GR gravity
(Deffayet & Menou 2007)
 - delay between arrival time of photons and gravitons:
improved limits on graviton mass ($\gamma m_0 c^2 = hf$; Kocsis et al. 2008)
 - frequency-dependence in delay: test Lorentz invariance
- Revolution for astronomy and astrophysics:
 - *accretion physics*: Eddington ratio and spectrum, as functions of BH mass and spin, orbital parameters
 - *quasar/galaxy co-evolution*: long-standing problem

Can we find EM counterparts?

- Sky position error from LISA is poor ($\sim 0.3 \text{ deg}^2$)
 - $10^{4-5} \rightarrow 10^{2-3}$ galaxies with LISA redshift info (i.e.: 3D)
 - perhaps a unique near-Eddington quasar (Kocsis et al. 2005)
- EM signature produced by merger is not understood
 - hard problem, requires gas physics + GR + radiation
- But ‘last parsec problem’ suggests gas needed
 - without gas, orbital decay / angular momentum loss time-scale exceeds Hubble time at $r \sim 1 \text{ pc}$
(Begelman, Blandford, Rees 1984)
- IF gas is still present at the time of GW-emitting phase
 - accretion onto one or both holes (or to post-merger binary)
 - modulations on orbital time-scale? post-merger shocks?
(Kocsis et al. 2006; 2008)

A Unique Quasar Counterpart?

TABLE 1
LISA MEASUREMENT ERRORS

	$\delta\mathcal{M}/\mathcal{M}$	$\delta\mu/\mu$	$\delta d_L/d_L$	$\delta\Omega$
best	0.8×10^{-5}	2×10^{-5}	2×10^{-3}	0.01 deg^2
typical	2×10^{-5}	9×10^{-5}	4×10^{-3}	0.3 deg^2
worst	0.8×10^{-3}	0.1	2×10^{-2}	3 deg^2

NOTE. — Assumed SMBH binary parameters: $m_1 = m_2 = 10^6 M_\odot$ and $z = 1$.

Vecchio (2004)

Angular and distance localization from GW signal alone depends on physical and orbital parameters and orientation

Kocsis, Frei, Haiman & Menou (2005)
Hughes & Holz (2005)

Angular Error: large LISA uncertainty (contains 10^{3-5} galaxies)

Distance Errors: - LISA $d_L(z)$ measurement

- Cosmological Model

- Peculiar velocity

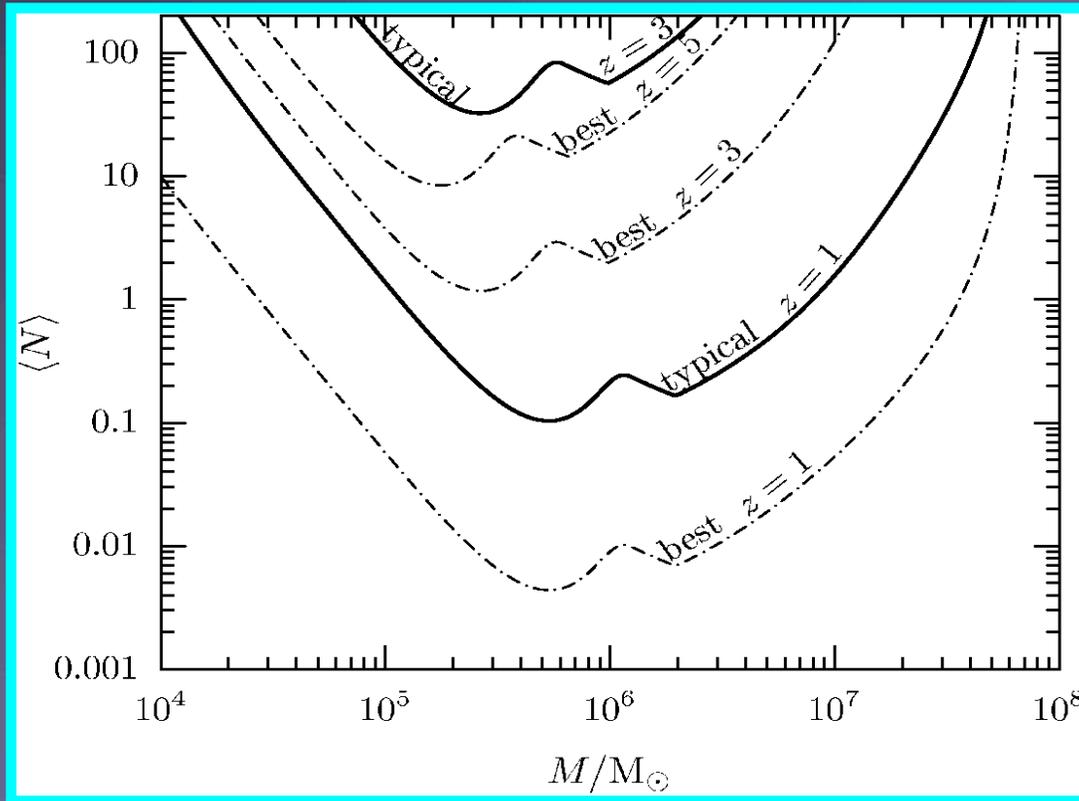
- Lensing-induced d_L -variations: $\Delta z \approx 0.002-0.03$

at $z \approx 0.3-1$



$\Delta z \lesssim 0.005$

Number of Quasars in 3D LISA Error Box



Assume:
SMBH coalescence
GW phase can be
accompanied by
luminous QSO
activity

Kocsis et al. (2005)

- Extrapolate known optical QSO LF to $M_{\text{BH}} \lesssim 3 \times 10^7 M_{\odot}$
- Assume $L/L(\text{edd}) \sim 0.3$, consistent with recent obs+models
- Compute mean number in error box (20% lensing correction)
- **Unique counterpart at $z < 1$ for $4 \times 10^5 M_{\odot} \lesssim M_{\text{BH}} \lesssim 10^7 M_{\odot}$**
- Can be extended to $z=3$ if BHs spin rapidly

Identify the Counterpart from Variability

AFTER THE MERGER IS COMPLETE:

- (1) Gravitational recoil at coalescence can cause strong shocks in circumbinary gas. Monitor 3D LISA error box ~months after the merger and look for prompt transient “afterglow”

(Lippai, Frei & Haiman 2008; Corrales, MacFadyen & Haiman 2008)

BEFORE AND DURING COALESCENCE:

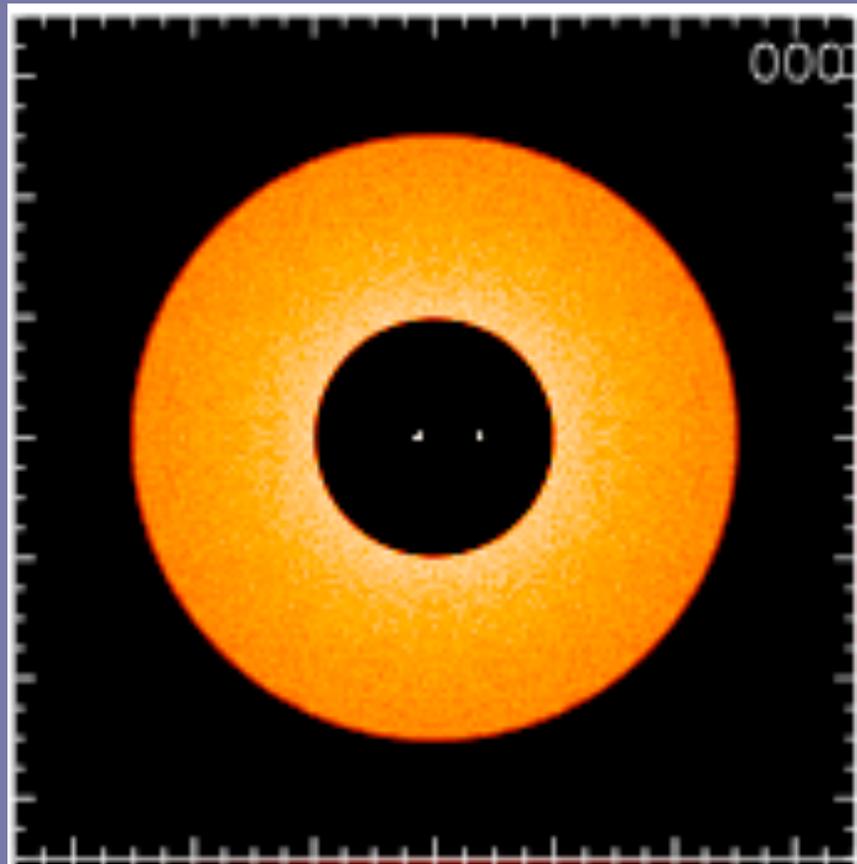
- (2) Can real-time LISA data-stream localize the source ~month **in advance**, so that a world-wide search can be triggered for **periodic variability** on the orbital timescale?

(Kocsis, Haiman, Menou & Frei 2007; Kocsis, Haiman & Menou 2008)

Cartoon Model of Binary+Gas Evolution

- Gas cools and settles into a thin circumbinary disk (Barnes 2002)
- Disk aligned with binary orbital plane (Bardeen & Peterson 1975)
(Ivanov et al. 1999)
- Torques from binary evacuate central cavity $r \sim 2a$
(Artymowicz & Lubov 1994)
- Binary orbit decays due to gas viscosity, cavity follows
- t_{GW} becomes shorter than t_{vis} when $r \sim \text{few } 100 R_S$
- Soon afterwards, disk ‘decouples’, cavity cannot follow at $r \lesssim 100 R_S$
- rapid GW-driven coalescence leaves ‘punctured disk’
(Milosavljevic & Phinney 2005)

Punctured disk



Cuadra et al. (2008)

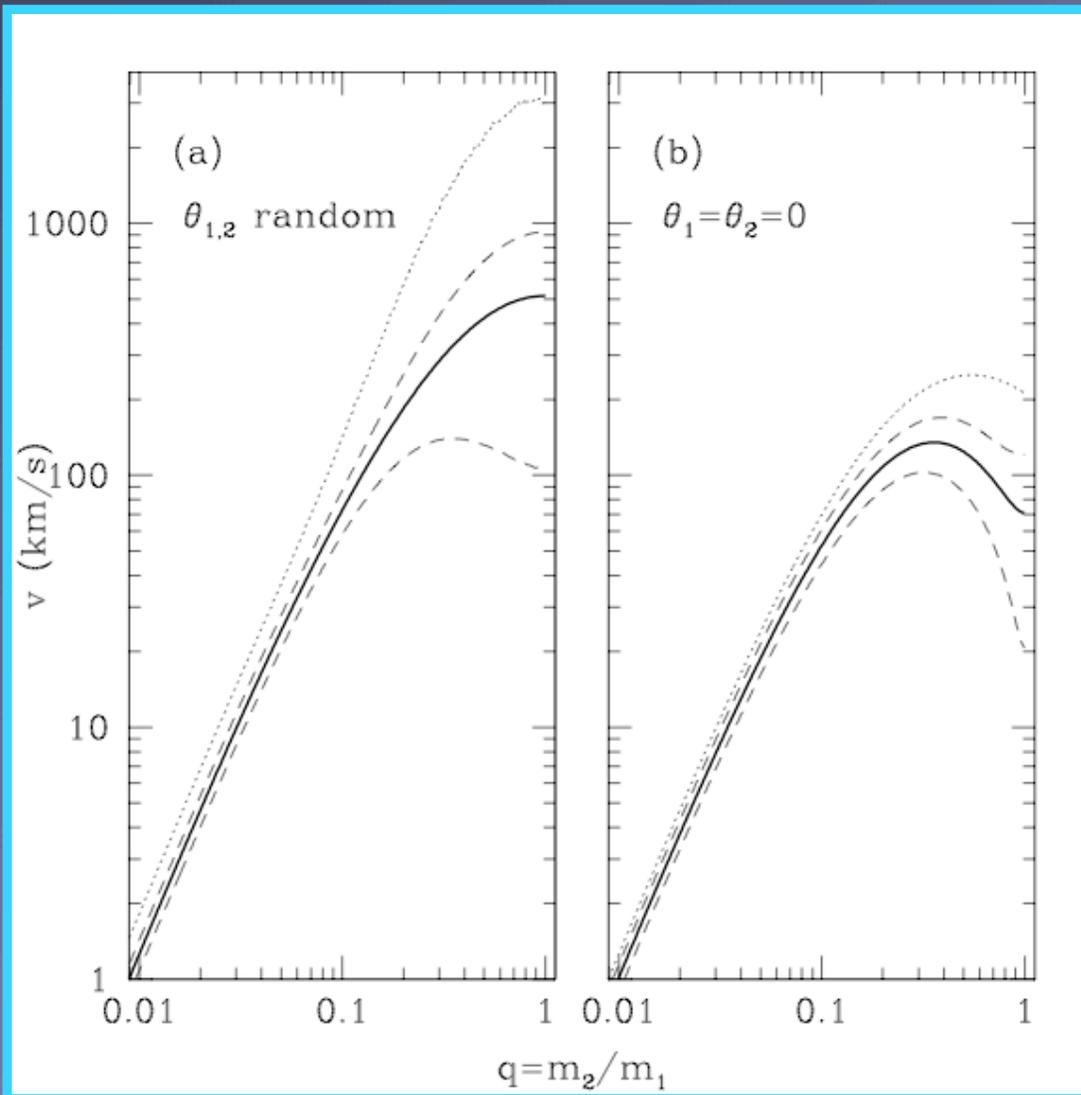
Gravitational Recoil

- **Gravitational radiation produces sudden recoil**
 - from conservation of linear momentum, near ISCO
 - kick velocity depends on mass ratio and on spin vectors
 - typical $v(\text{kick}) \sim \text{few} \times 100 \text{ km/s}$ (Baker et al. 2006, 2007)
 - maximum $v(\text{kick}) \sim 3\text{-}4,000 \text{ km/s}$ (Gonzalez et al. 2007)
 - directed in the plane if spins aligned, generally out of the plane otherwise

- **What is the response of the circumbinary disk?**
 - can we expect prompt EM signal, within years, so that it is useful for selection among LISA candidates?
(Lippai et al. 2008)

Kick Velocity Distribution

Tanaka & Haiman (2008) from Baker et al. (2008)



10^6 realizations
with random
spin magnitude
in the range
 $0 < a_1, a_2 < 0.9$

($\pm 1\sigma$ range shown)

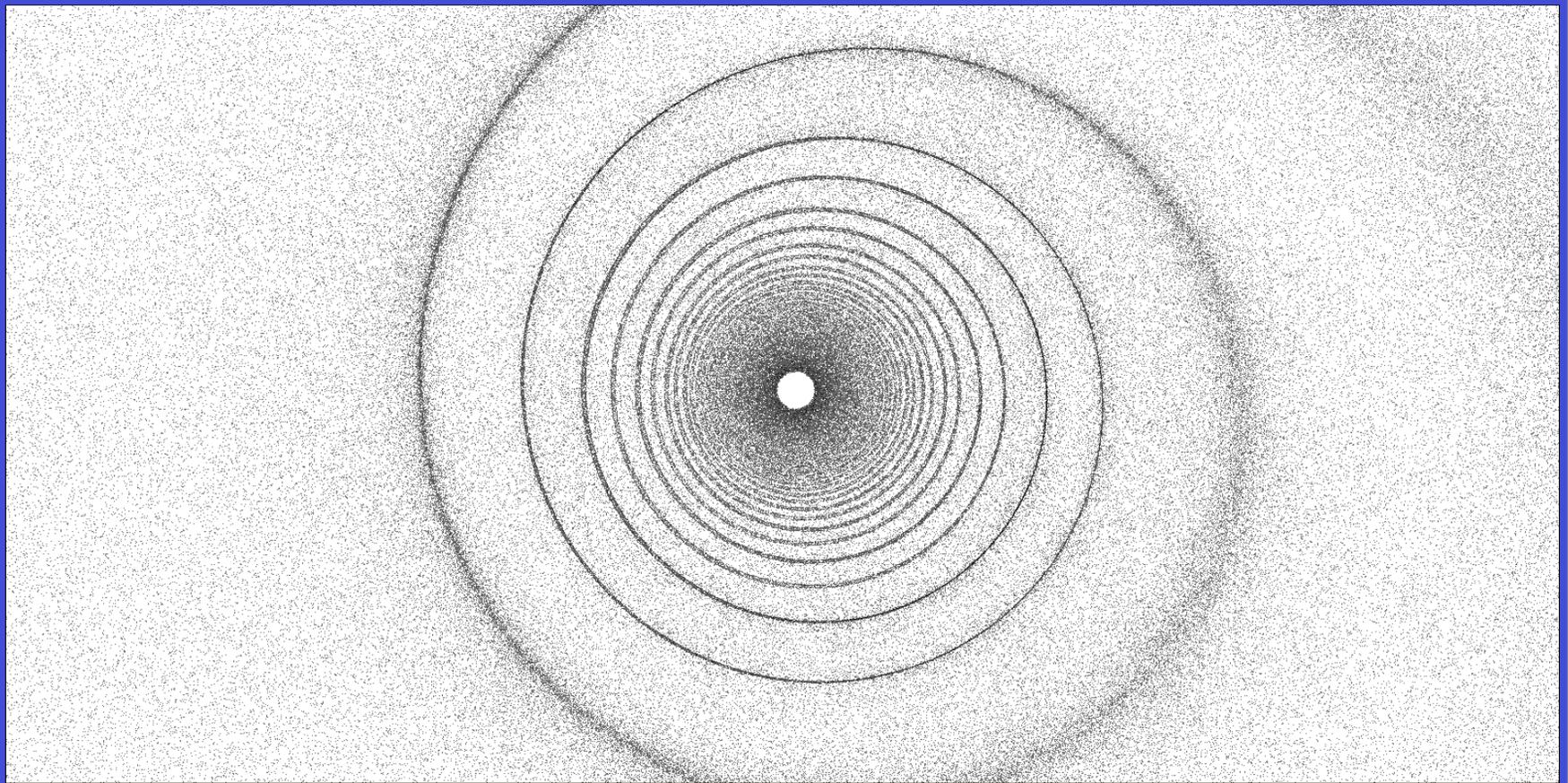
Effect of Kick on Circumbinary Disk

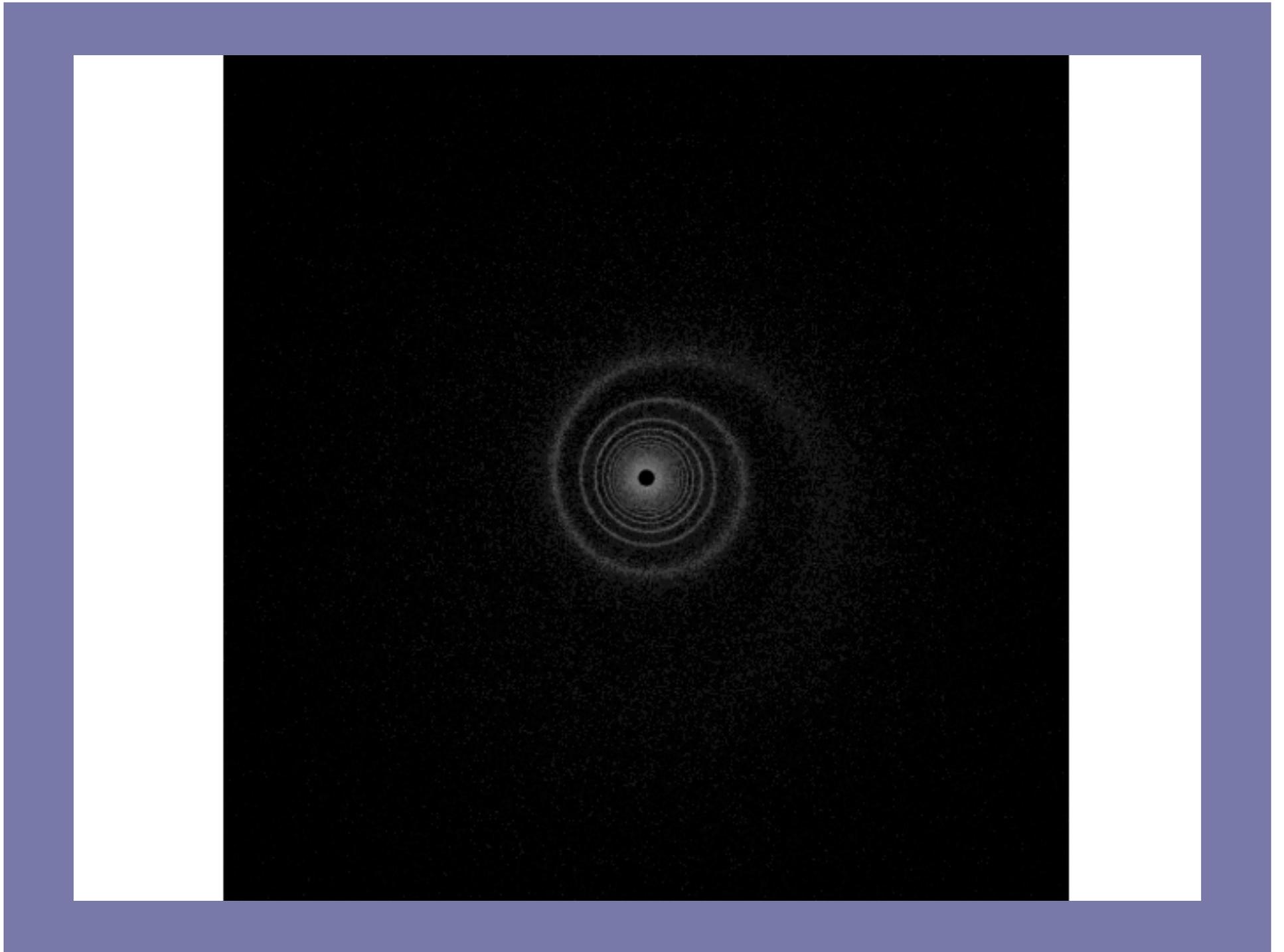
Lippai, Frei & Haiman (ApJL 2008)

- **Properties of disk:**
 - geometrically thin (cold) accretion disk, susceptible to shocks
 - inner cavity, evacuated by torques (out to $\sim 100 R_s$)
 - disk gravitationally unstable beyond $\sim 10,000 R_s$
 - $v(\text{orbit}) \sim 20,000 \text{ km/s} \rightarrow 2,000 \text{ km/s}$
 - inner disk tightly bound to binary, outer disk weakly bound
 - disk mass low ($M_{\text{disk}} \sim 10^{-4} M_{\text{BH}}$): no effect on BH trajectory
- **Response of pressureless (“dark matter”) disk:**
 - start with massless test particles on circular orbits
 - add instantaneous $v(\text{kick})$, parallel or perpendicular to disk
 - follow Kepler orbits (ellipses) for $N=10^6$ particles

Planar Kick Results in a Spiral Caustic

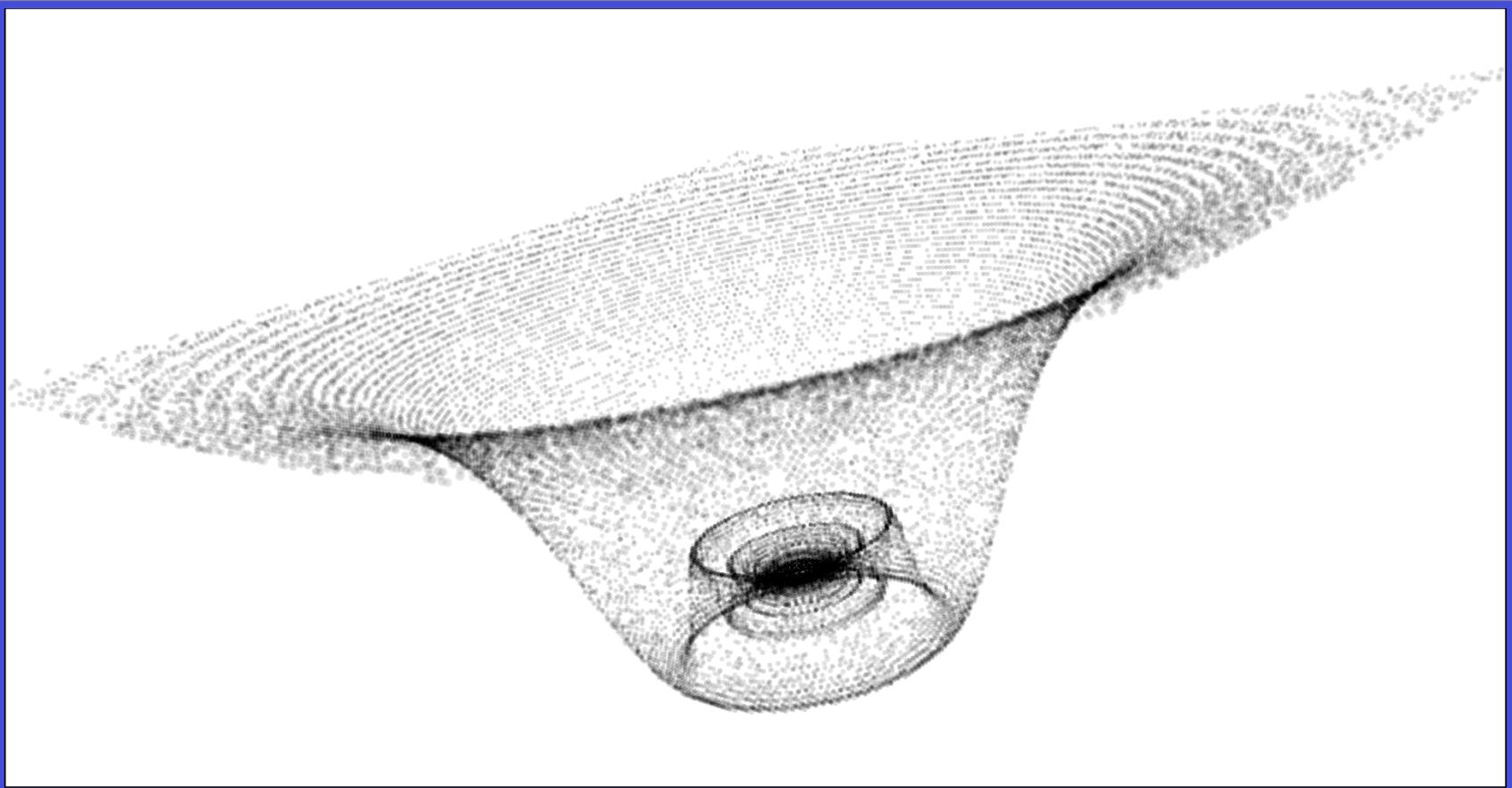
$$M_{\text{BH}} = M_1 + M_2 = 10^6 M_{\odot} \quad (R_{\text{cavity}} = 100 R_s = 2 \text{ AU})$$
$$v_{\text{kick}} = 500 \text{ km/s} \quad (\text{kick in the disk plane})$$
$$t = 90 \text{ days} \quad (t_{\text{cavity}} = R_{\text{cavity}} / v_{\text{kick}} = 7 \text{ days})$$





Perpendicular kick: Concentric Density Enhancements

(otherwise same parameters)



Expected Caustic Properties

Consider caustic formed from material with annulus $\Delta R \ll R$
and use epicyclic approximation:

epicyclic amplitude: $\Delta R \sim (v_{\text{kick}}/v_{\text{orbit}}) \times R$

caustic forms at time: $t \sim [(d\Omega/dR) \times \Delta R]^{-1}$

$$\rightarrow t \sim [(d\Omega/dR) \times (v_{\text{kick}}/v_{\text{orbit}}) \times R]^{-1}$$

use $d\Omega/dR \propto \Omega/R$

$$\rightarrow t \sim [\Omega (v_{\text{kick}}/v_{\text{orbit}})]^{-1} = R/v_{\text{kick}}$$

propagation speed: $R/t = v_{\text{kick}}$

infall speed: $v_{\text{shock}} \sim \Delta R/t \sim \Delta R/(R/v_{\text{kick}}) \sim v_{\text{kick}}^2/v_{\text{orbit}}$

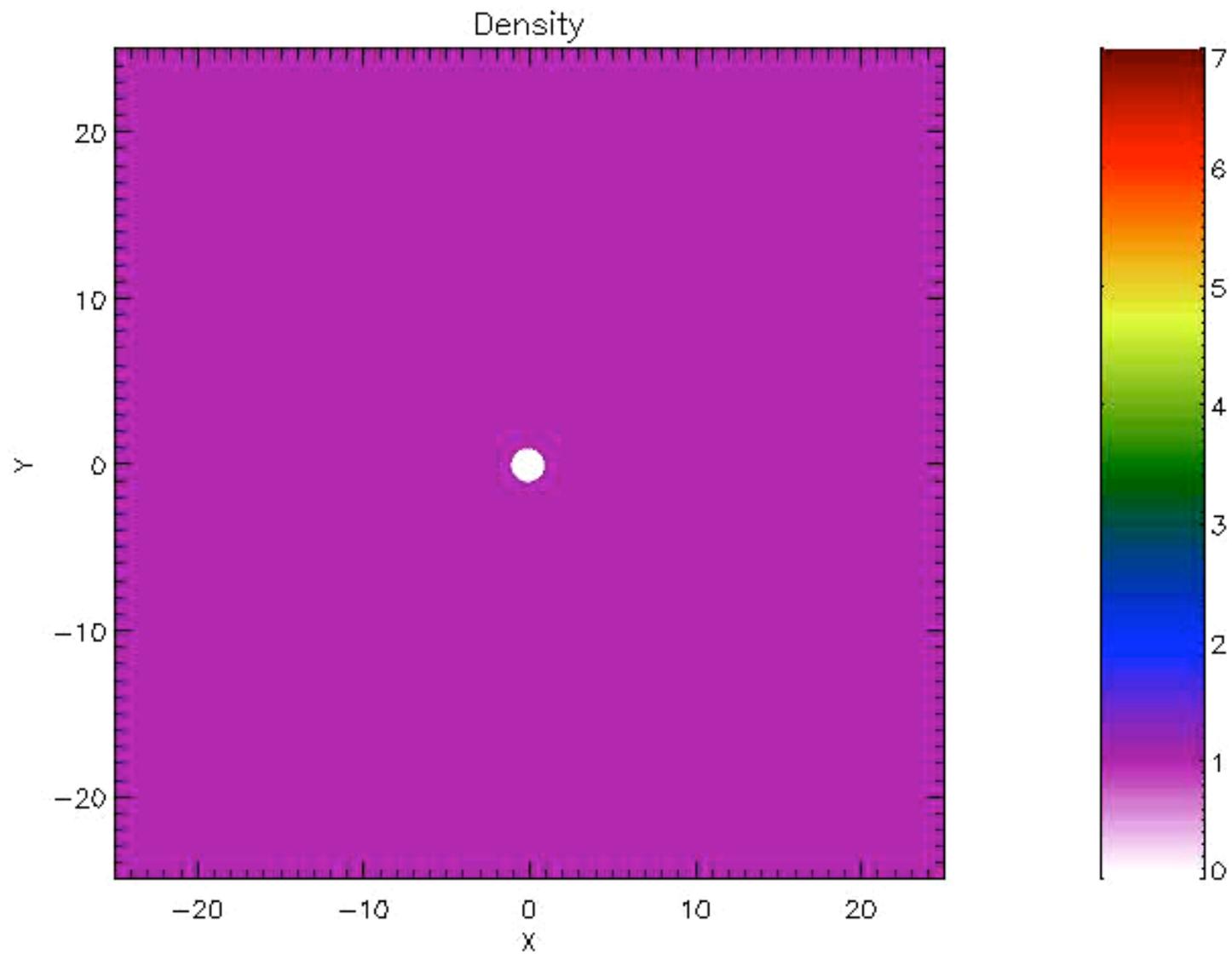
Implications of prompt spiral caustics

- Suggests prompt “afterglow” for SMBH coalescence:
 - caustic propagates outward with speed $\sim v_{\text{kick}}$
 - infall speed into caustic is $v_{\text{caustic}} \sim v_{\text{kick}}^2 / v_{\text{orbit}}$
 - v_{caustic} becomes supersonic beyond $\sim 700 R_s$ (at > 25 km/s)
 - gas shocks may produce strong emission (at > 50 days)
- Can speculate about properties of afterglow:
 - shocked gas heated to $v_{\text{shock}} \sim v_{\text{caustic}} \sim 25 - 80$ km/s
 - $L_{\text{disk}} \sim 1/2 M_{\text{disk}} v_{\text{shock}}^2 / t_{\text{shock}}$
 - $M_{\text{disk}} \sim 50 - 1,200 M_{\odot}$ $t_{\text{shock}} \sim 50$ days - 2 years
 - $L_{\text{disk}} \sim 6 \times 10^{-4} - 2 \times 10^{-2} L_{\text{edd}}$ not negligible.
 - Hardens from UV to soft X-ray (opposite of GRB afterglow)

Impact of Gas Dynamics

Corrales, MacFadyen & Haiman (2009)

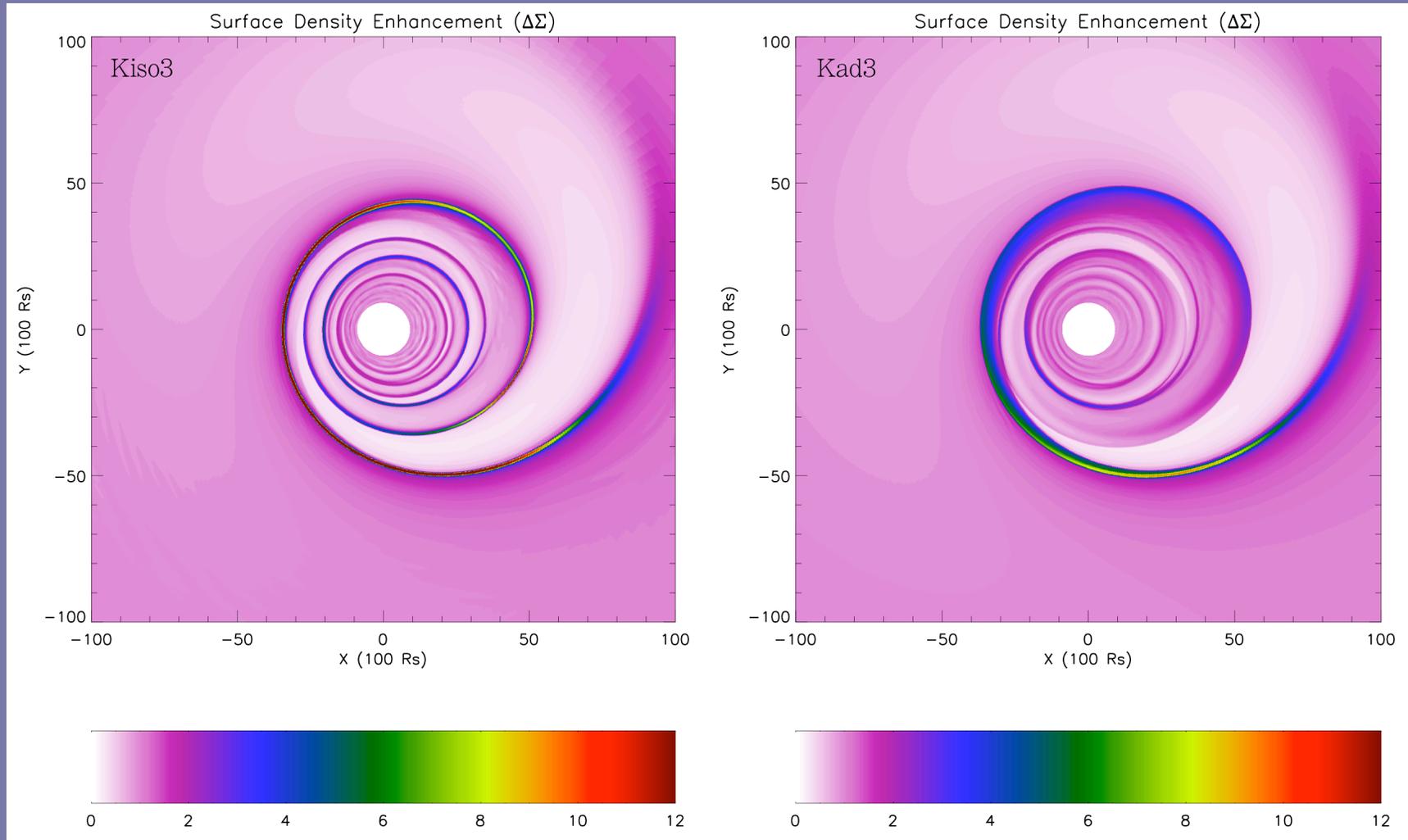
- Sudden ‘shaking’ of disk launches prompt sound waves
- Sound waves can steepen into shocks
- Hydro simulation
 - adaptive mesh refinement (AMR) code FLASH
 - $v_{\text{kick}} = 500 \text{ km/s}$
 - equation of state: isothermal or adiabatic
 - vary temperature: $5000 - 5 \times 10^5 \text{ K}$



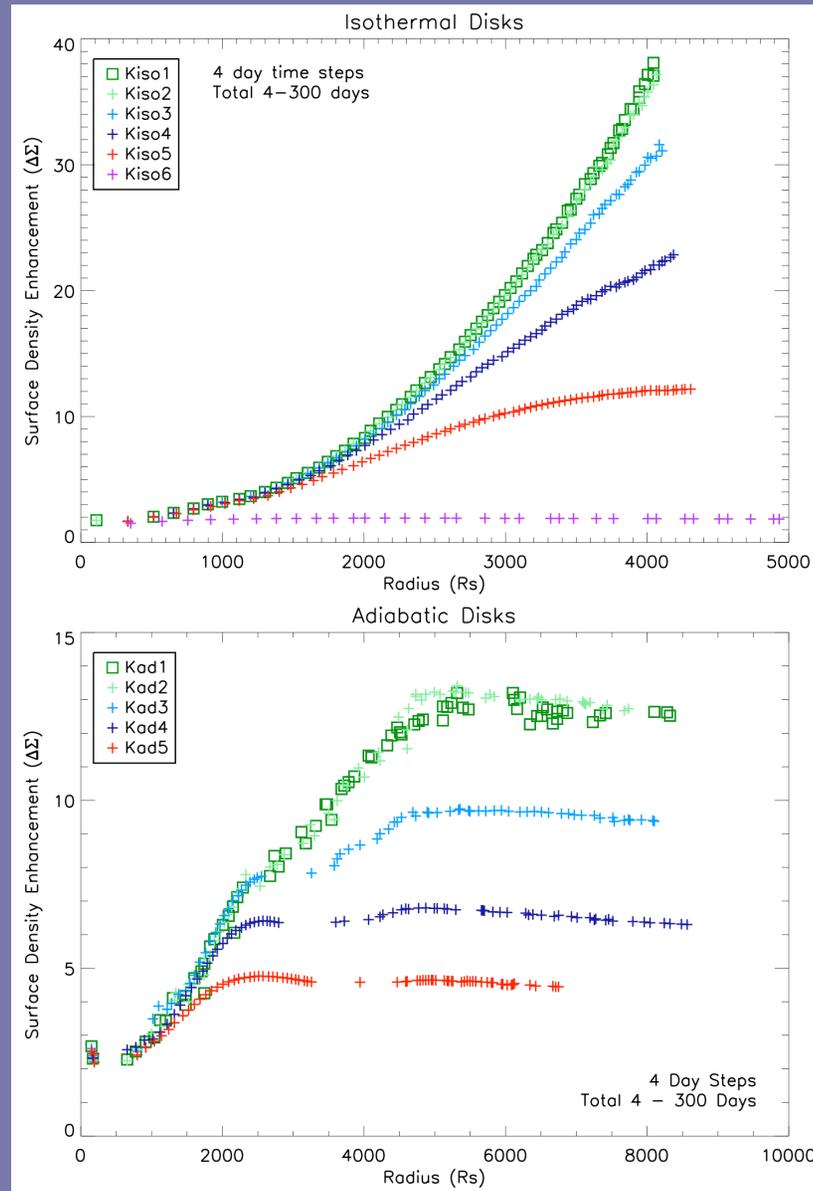
time = 0.000 ps
number of blocks = 7808
AMR levels = 5

Disk Surface Density

Corrales, MacFadyen & Haiman (2009)



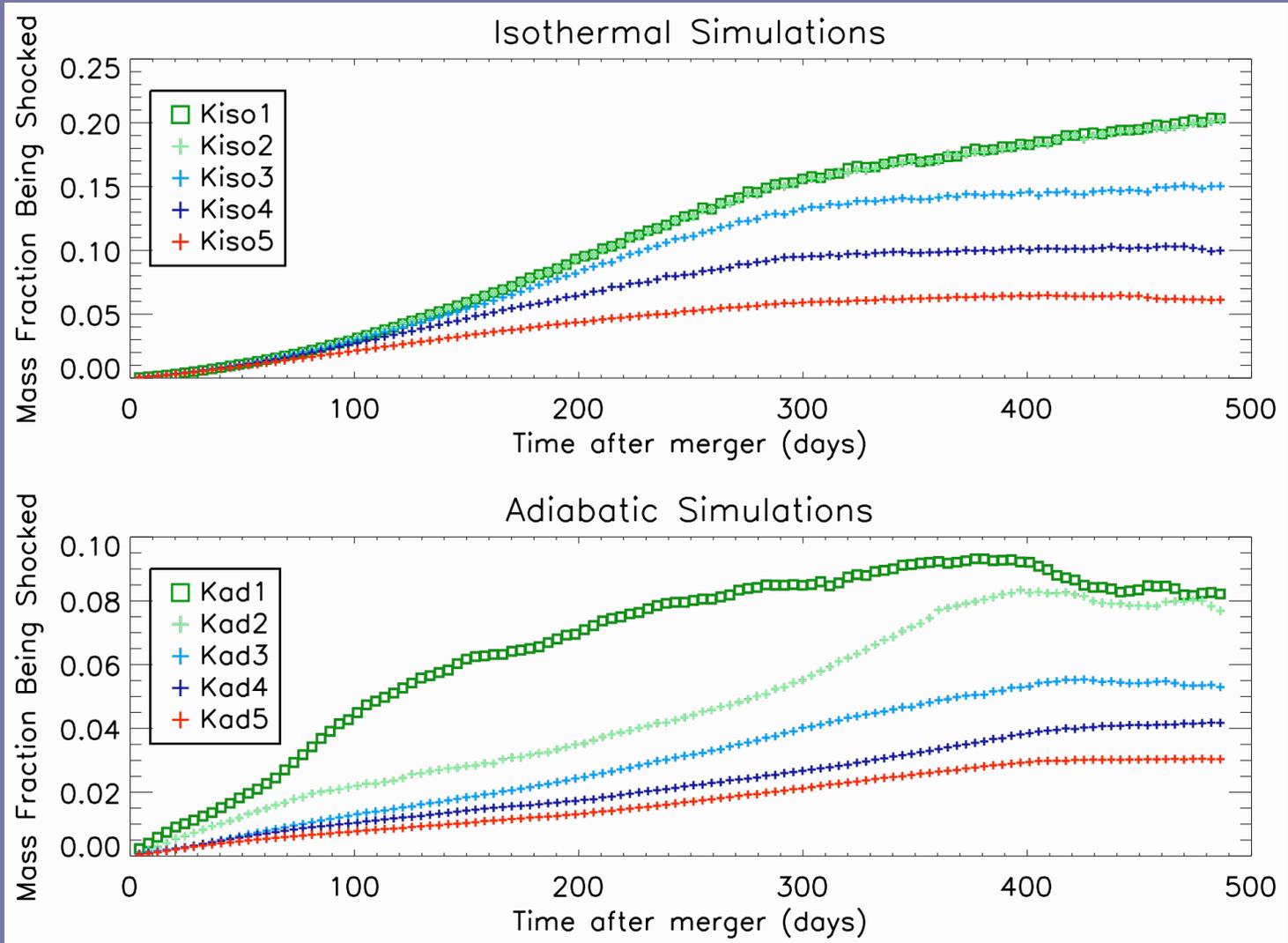
Over-density in Spiral Shocks



Corrales,
MacFadyen
Haiman (2009)

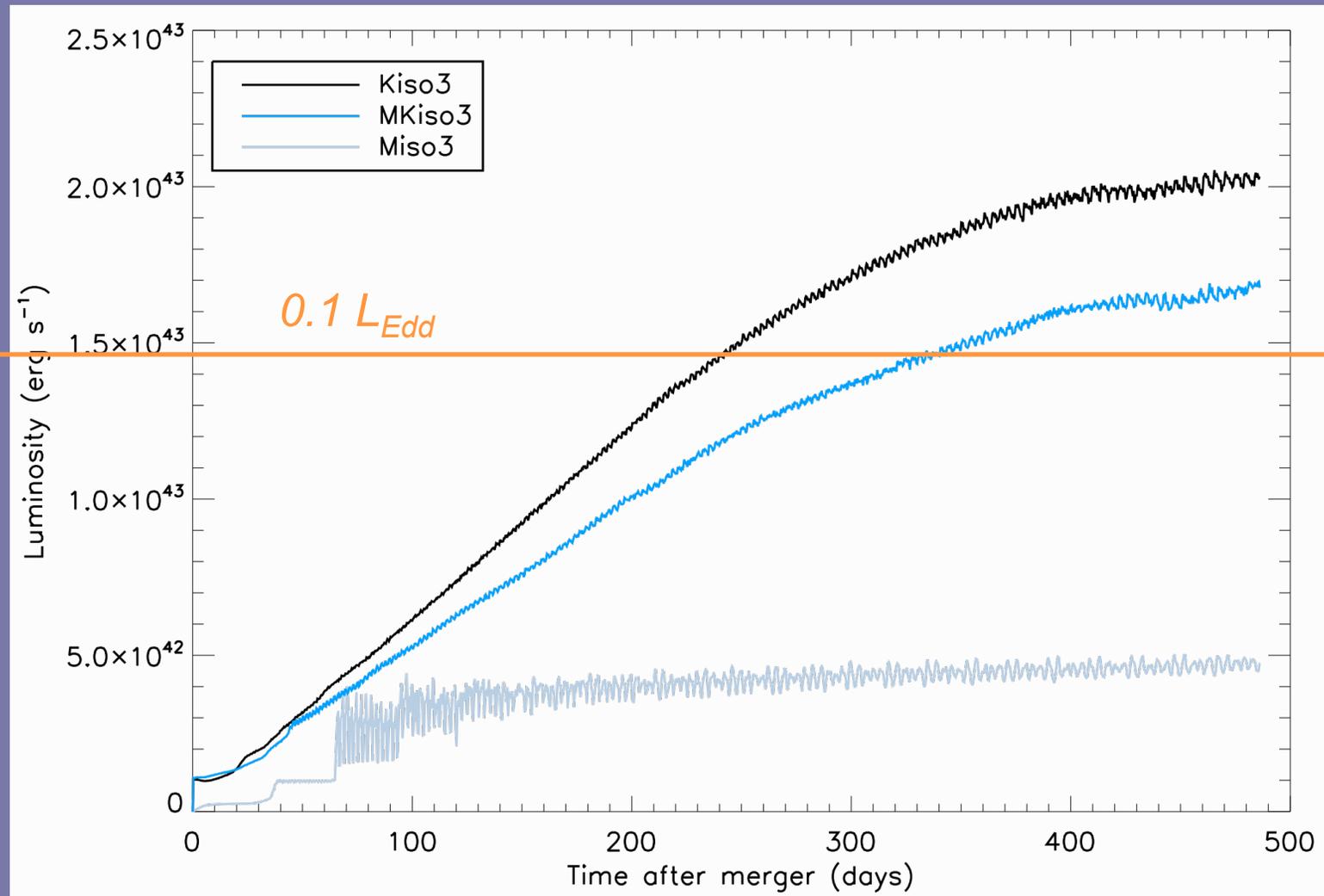
Shocked gas fraction

Corrales, MacFadyen & Haiman (2009)



Light Curve

Corrales, MacFadyen & Haiman (2009)



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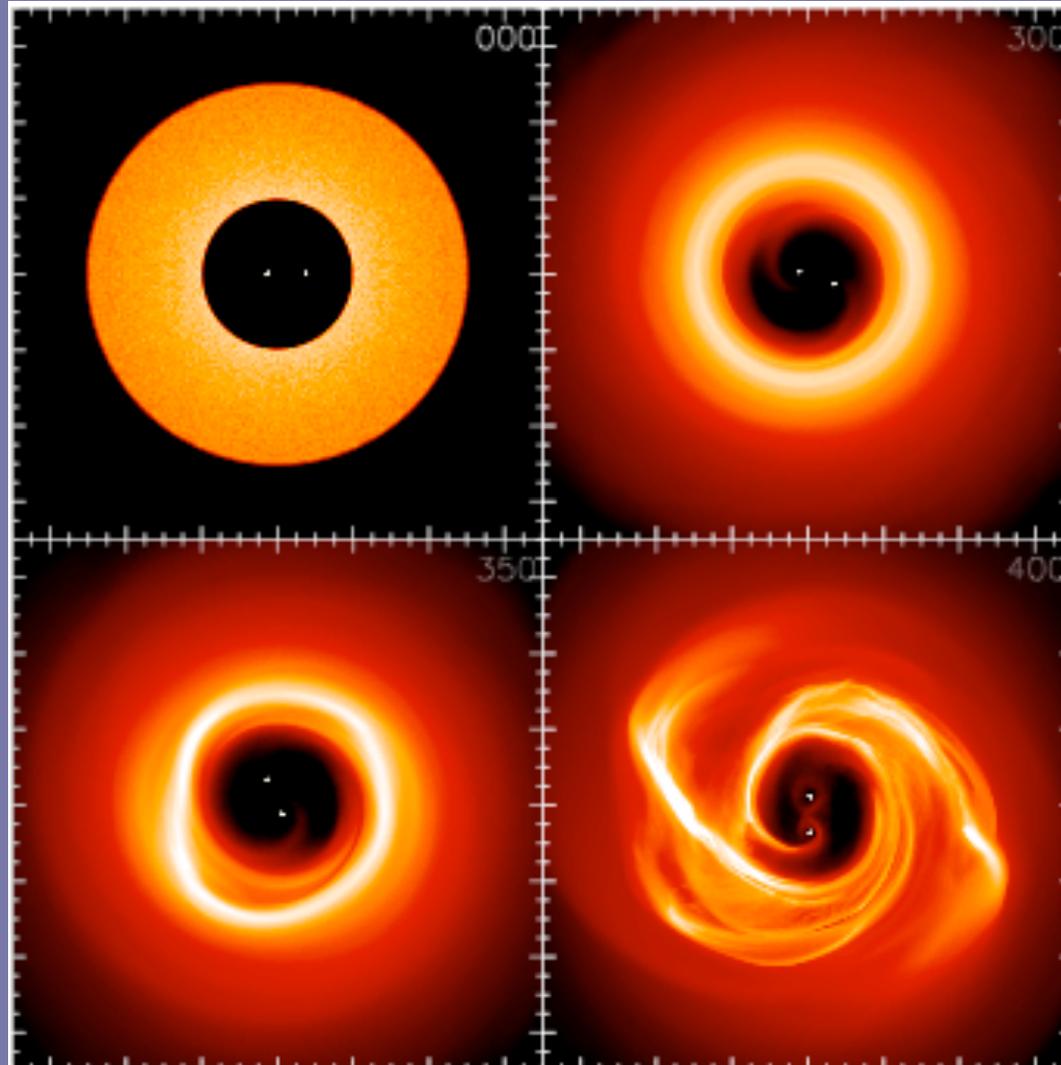
BEFORE AND DURING COALESCENCE:

- (2) Can real-time LISA data-stream localize the source ~month **in advance**, so that a word-wide search can be triggered for **periodic variability** on the orbital timescale?

(Kocsis, Haiman, Menou & Frei 2007; Kocsis, Haiman & Menou 2008)

Gas Near BHs Prior to Merger

Cuadra et al. (2008)



Localizing a LISA source

What is the sky position error (and shape) in last weeks of merger?

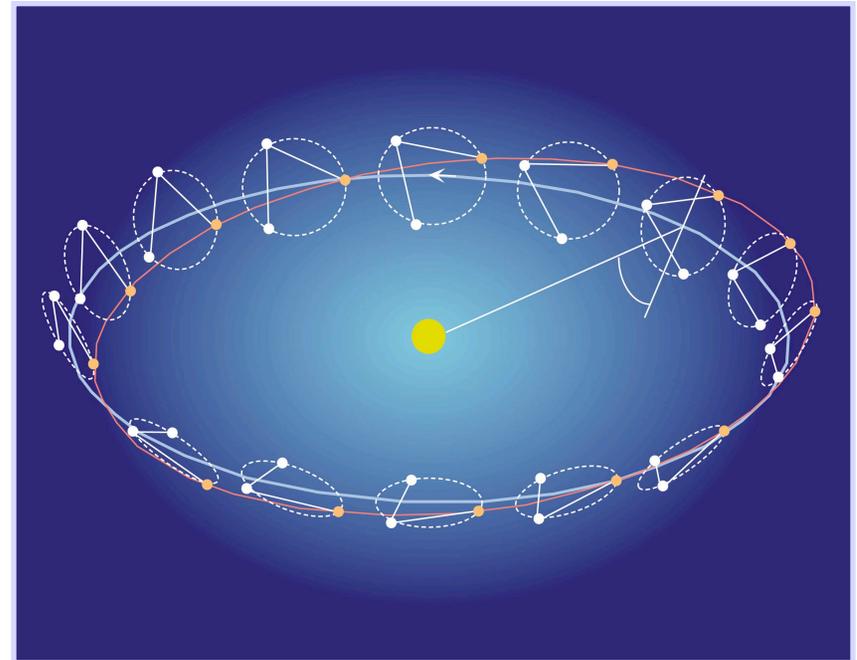
* Orbital modulation of the GW signal

$$h(t) = \frac{G(\Omega, \dot{\Omega}, t)}{d_L} h_0(t, \mathcal{M}_z, \mu_z, \text{spins}) + \text{noise}(t)$$

Detector orientation: (Ξ_D, Φ_D)

Source direction: (θ_N, ϕ_N)

Source orientation: (θ_L, ϕ_L)



* Amplitude modulation is periodic with

$$f_{\oplus} \equiv \frac{1}{\text{yr}} \ll f_{\text{GW}}(t)$$

Harmonic Mode Decomposition

(Kocsis, Haiman & Menou 2007)

- Measured GW signal can be written in an equivalent form

$$h^{I,II}(p;t) = h_c(p_{\text{fast}}, p_{\text{spin}}; t) \times h_m^{I,II}(p_{\text{slow}}; t)$$

- Parameters dependence decoupled in three groups:

$$p_{\text{slow}} \equiv \{d_L, \Omega\},$$

$$p_{\text{fast}} \equiv \{\mathcal{M}_z, \mu_z, t_{\text{merger}}, \phi_{\text{ISCO}}\},$$

$$p_{\text{spin}} \equiv \{2 \text{ spin magnitudes}, 4 \text{ spin angles}\}$$

- The angular piece can be further simplified

$$h_m(p_{\text{slow}}(0), t) = d_L(z)^{-1} \sum_{j=-4}^4 g_j(p_{\text{slow}}(0)) e^{ij\omega_{\oplus} t}$$

**Time independent
angular dependence**

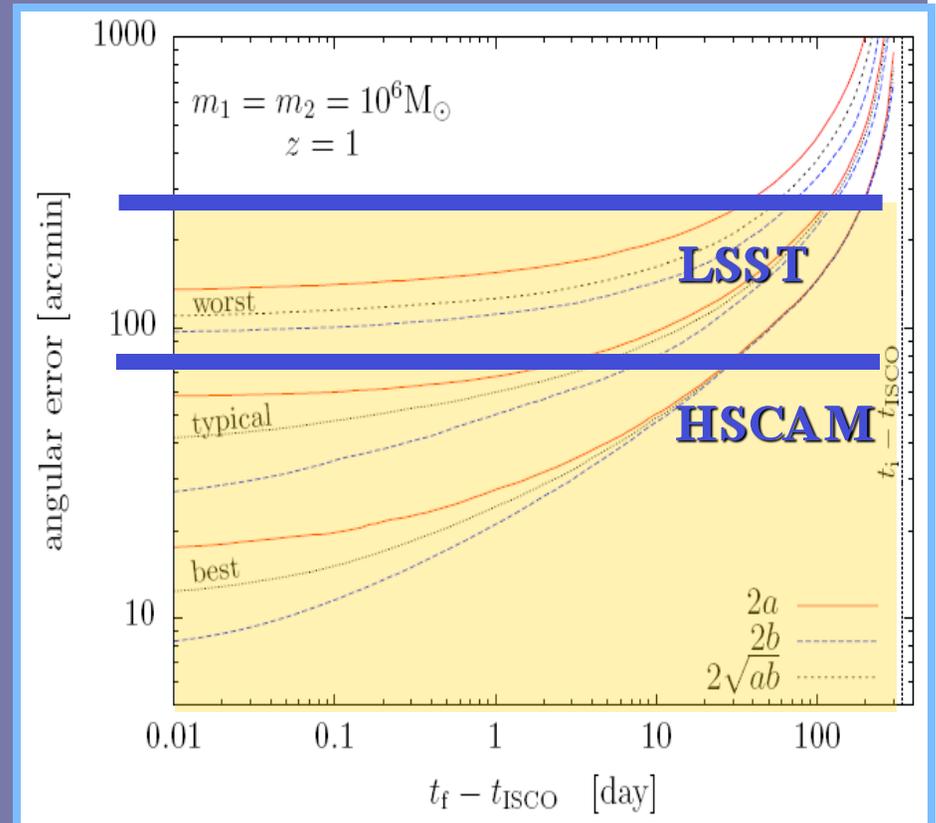
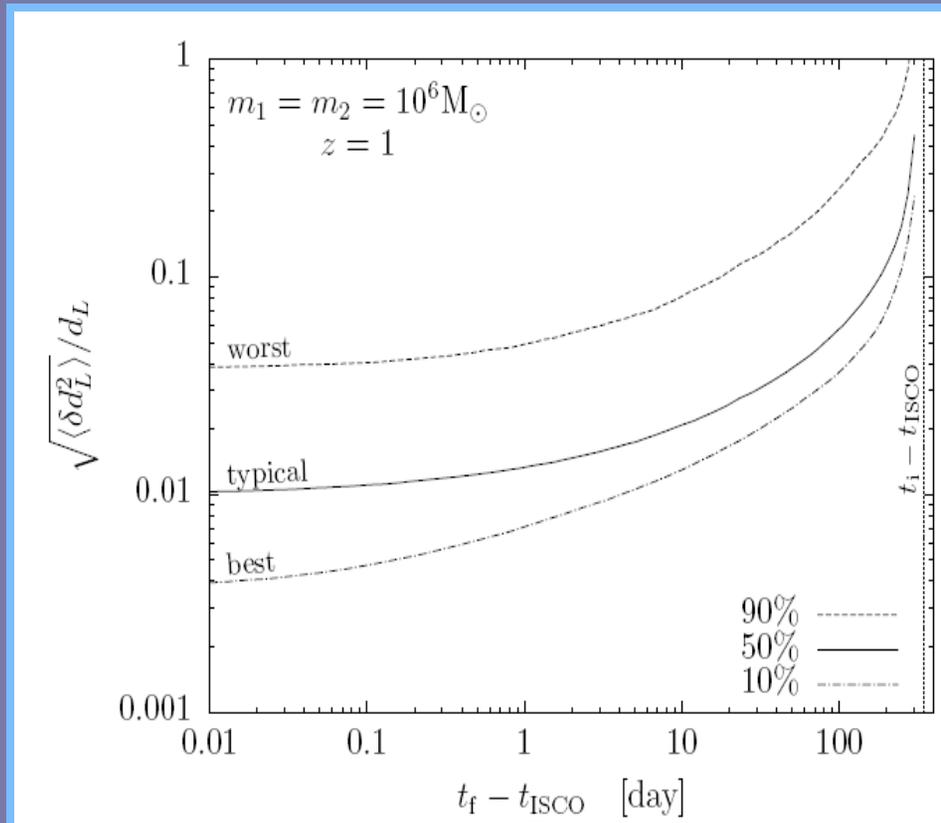
Orbital harmonics

Time dependence of localization

(HMD method + Fisher matrix: Kocsis, Haiman, Menou & Frei 2007)

distance uncertainty

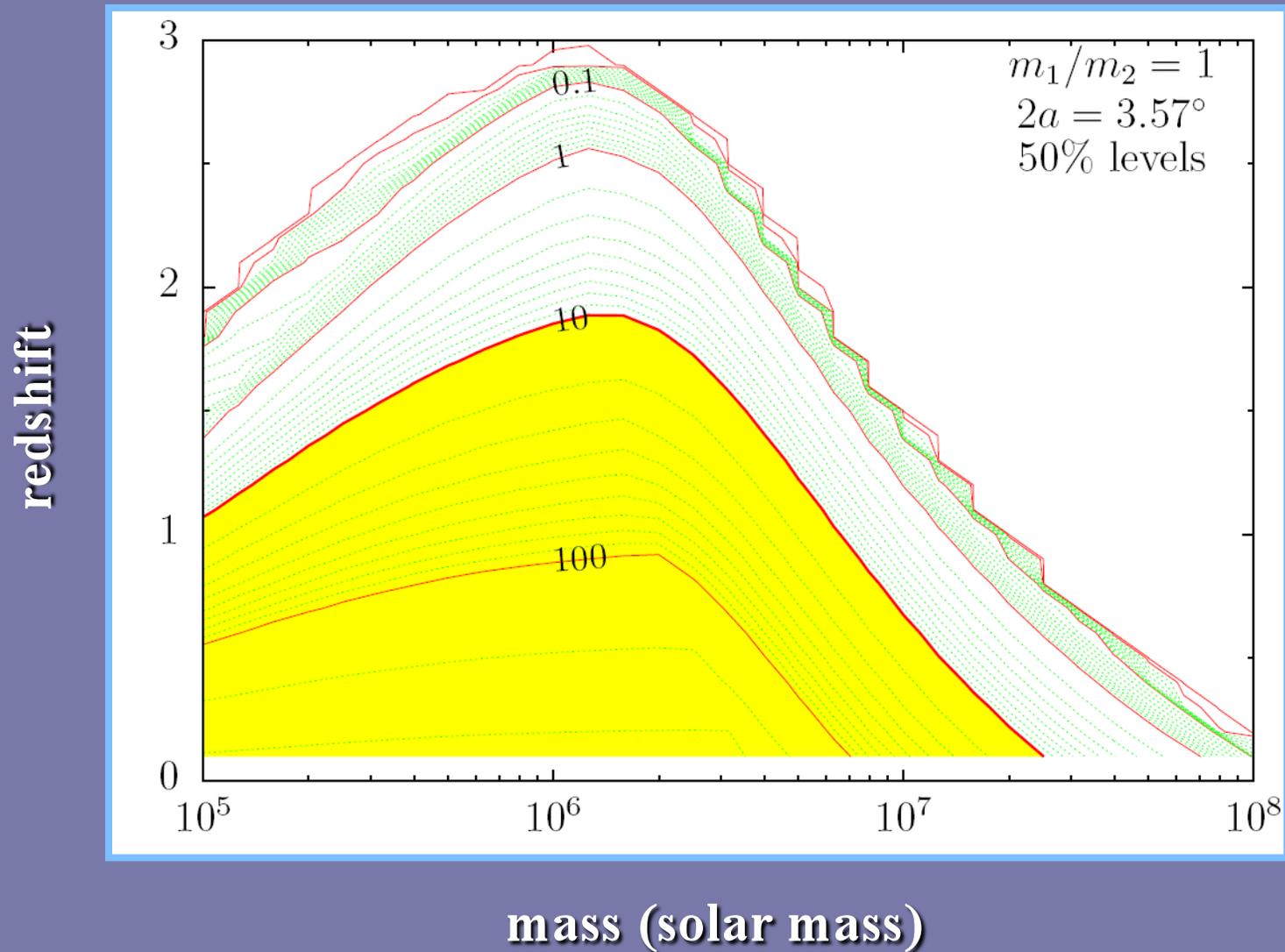
sky position uncertainty



Errors typically stop improving ~10 days before ISCO

How much advance notice?

Look-back time when sky position error shrinks down to $\sim 10 \text{ deg}^2$



A world-wide monitoring campaign with wide FOV instruments

(Kocsis, Haiman & Menou 2008)

GW source localized $\sim 2-3$ weeks before merger

- Monitor sources in few \times deg² field
- variability 24-27 mag on timescales of hours to minutes
(1-10% $L_{\text{Eddington}}$ for $M_{\text{BH}} = 10^{6-7} M_{\odot}$ at $z=1-2$)
- Correlate EM signal with GW template over 10^{2-3} cycles
- Sky position error shrinks to \sim ten arcmin in last few days
- Events with favorable geometry can be identified in advance

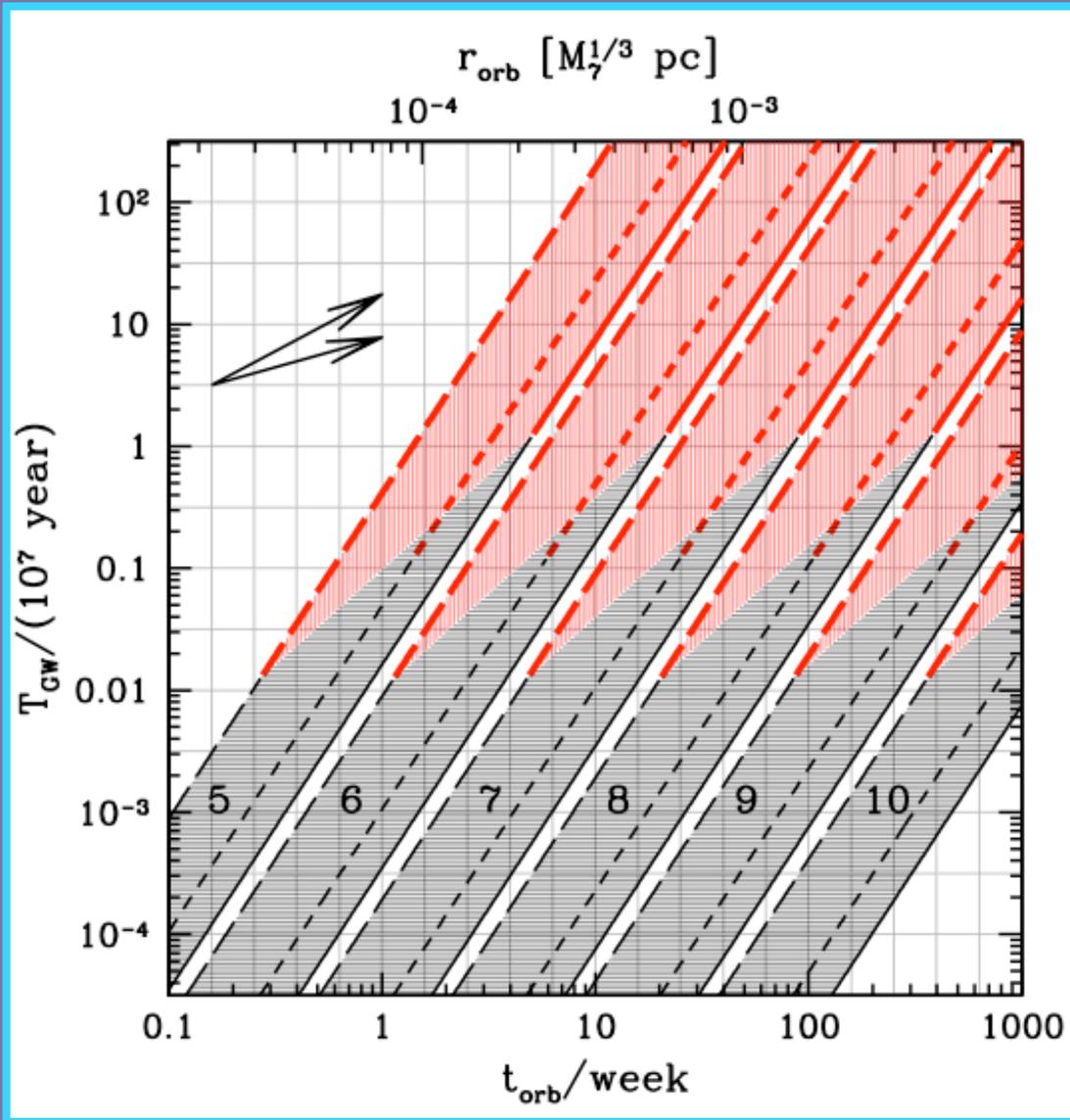
Do we have to wait for LISA?

(Haiman, Kocsis & Menou 2008)

- **OPTIMISTIC ASSUMPTIONS:**
 - binary is producing bright emission ($\sim 30\% L_{\text{edd}}$)
 - non-negligible fraction ($\sim 10\%$) of this emission is variable
 - clearly identifiable period $t_{\text{var}} \sim t_{\text{orbit}}$
 - orbital evolution driven by GWs below $r \lesssim 10^4 R_S$
 - one-to-one correspondance between BH mergers and quasars
- **CAN WE IDENTIFY SUCH GW-DRIVEN BINARIES ?**
 - GW-driven binary = periodically variable quasar
 - fraction of quasars with period t_{var} :

$$f_{\text{var}} = \frac{N_{\text{var}}}{N_{\text{tot}}} = \left(\frac{10^7 \text{ yr}}{t_Q} \right) \left(\frac{t_{\text{var}}}{19.8 \text{ weeks}} \right)^{8/3} \left(\frac{(1+z)M_{\text{tot}}}{10^6 M_{\odot}} \right)^{-5/3} \left(\frac{4q}{(1+q)^2} \right)^{-1}$$

Time spent at each orbital separation



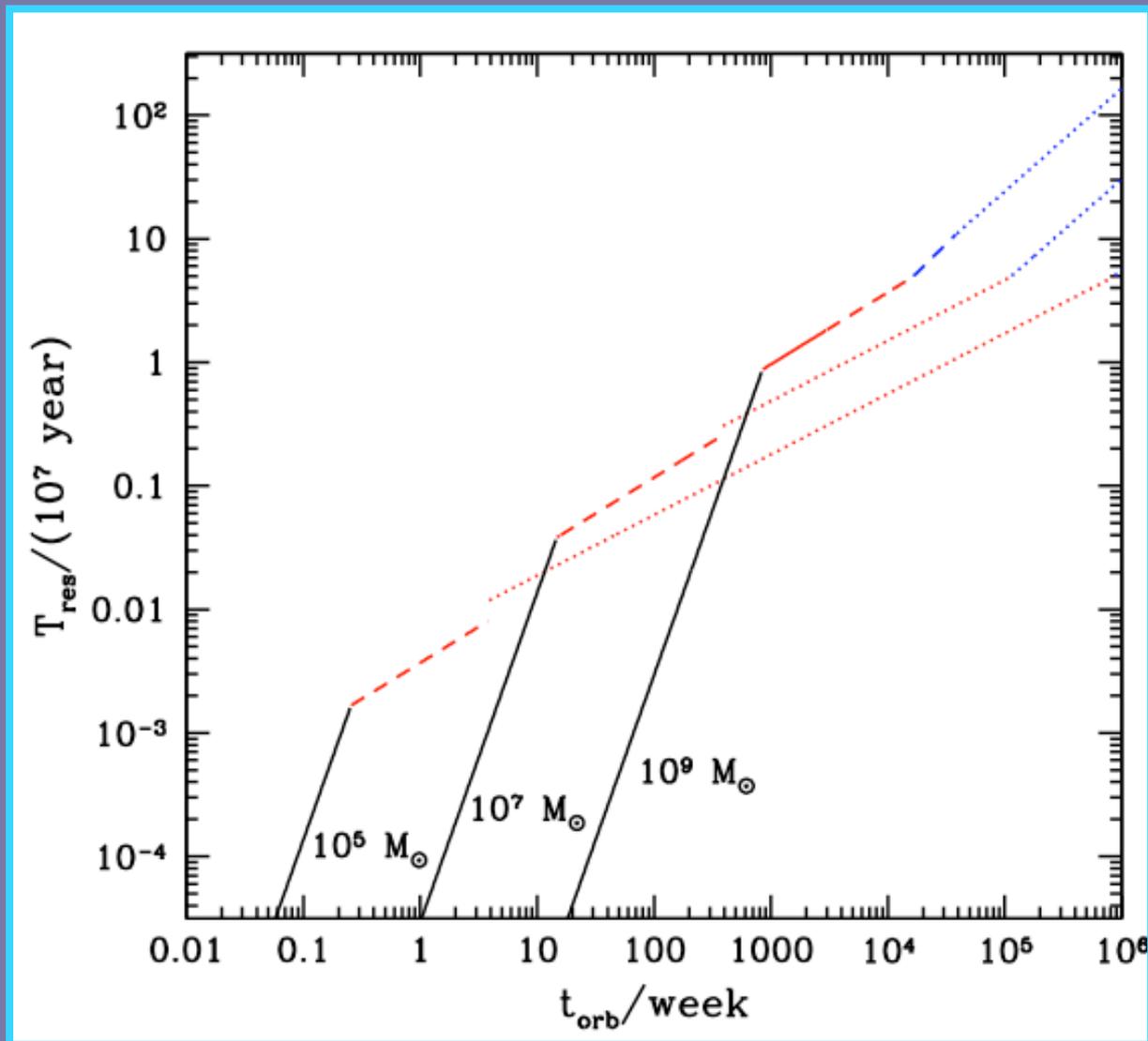
Assume:

- $t_{\text{mg}} = \text{time-to-merger}$ (Newtonian approx.)

Trade-off:

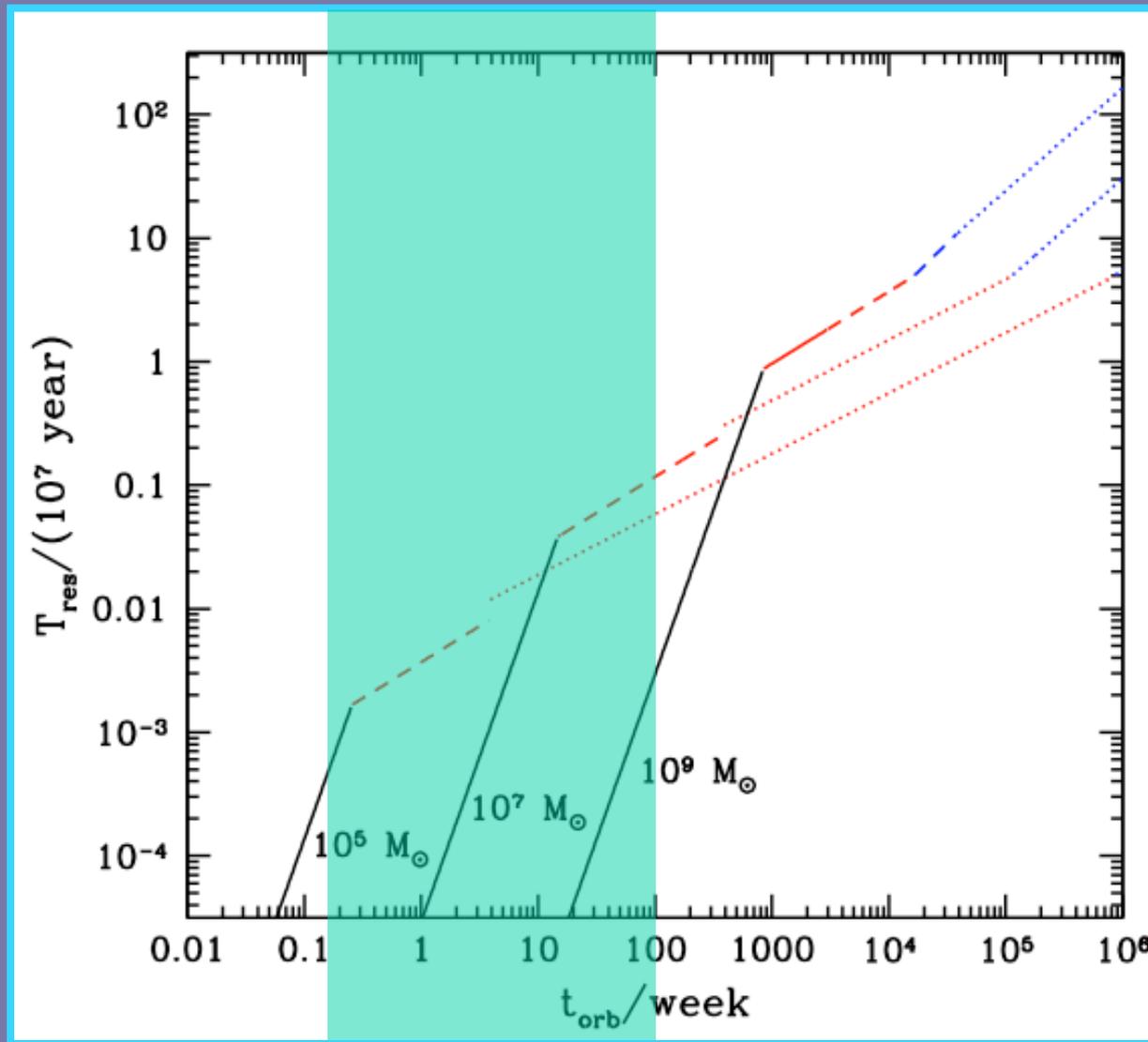
- periodicity more common among low-mass BHs, but they are faint
- high-mass BHs, are brighter, but periodic sources are more rare

Residence time: disk physics



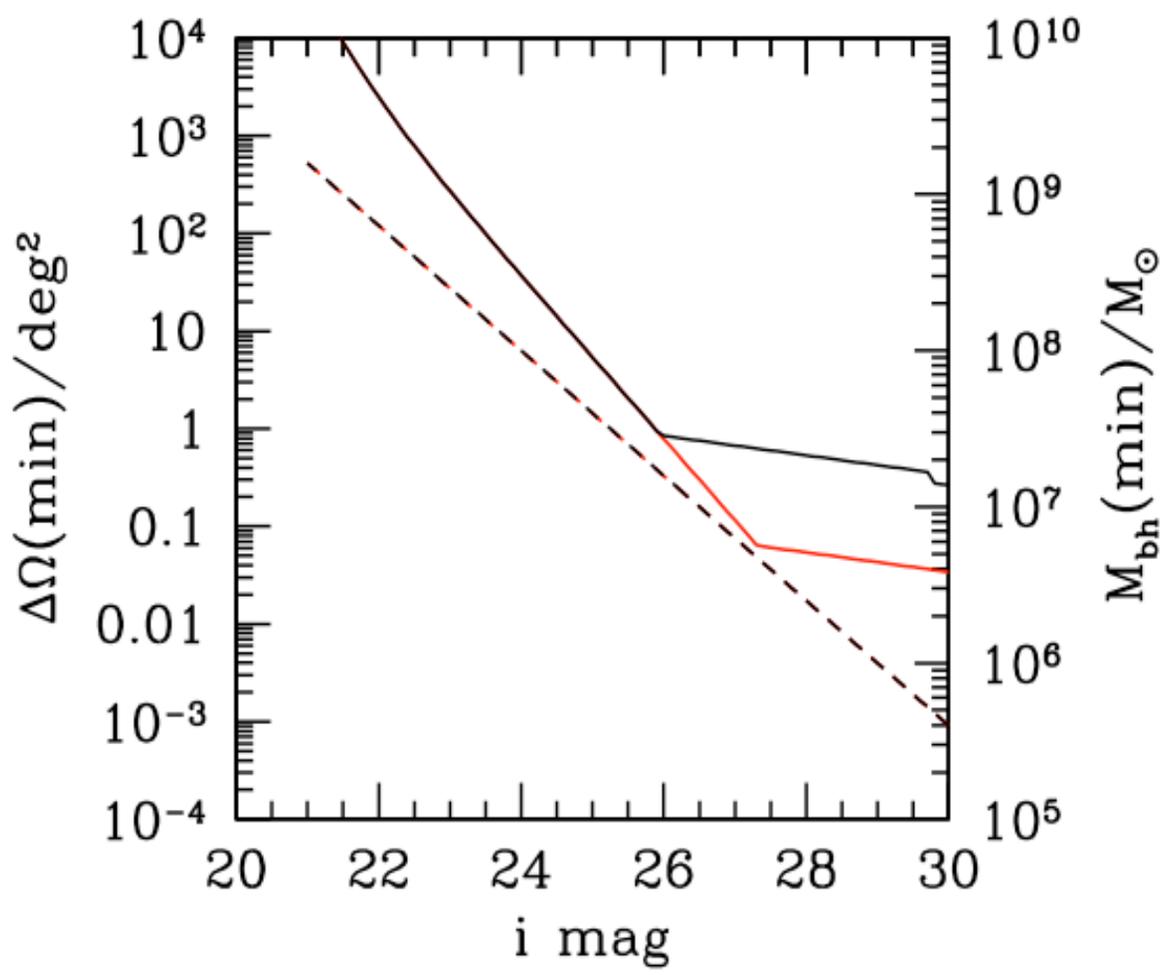
- “alpha disk”
- steady state
- equal-mass
- piece-wise power-law asymptotic solutions

Residence time: disk physics



- “alpha disk”
- steady state
- equal-mass
- asymptotic piece-wise power-law solutions

Requirements for an (optical) survey for finding periodic variable sources



Require:

- ≥ 5 sources @ $t_{\text{var}} \leq 20$ wk
- ≥ 100 sources @ $t_{\text{var}} \leq 1$ yr

Assume:

- $f_{\text{Edd}} = 0.3$
- $f_{\text{var}} = 0.1$
- $t_Q = 10^7$ yr
- Hopkins et al. QSOLF

Conclude:

- wide survey best to probe GW-decay
- disk physics at $i \sim 25$

Conclusions

1. Gravitational **recoil** launches prompt outward-moving **spiral shock wave** in circumbinary disk → produce a detectable transient afterglow (hardening with time?)
2. **Advance localization** possible **weeks-months** before merger, to within a **few square degrees**, triggering monitoring campaign with wide FOV telescopes
3. decaying binaries may be identifiable in a search for **periodic variability** among AGN, even before LISA, utilizing the scaling of occurrence rate (e.g. $f_{\text{var}} \propto t_{\text{var}}^{8/3}$)