

RRI
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Hydrodynamic Scales of Integrable Many-Particle Systems

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⇒ Calogero fluid ⇐

arXiv: 2301.08504

integrable many-particle systems

Lieb-Liniger δ -Bose gas (1963)

XYZ spin chains

Toda lattice (quantum + classical)

→ Calogero fluid ←

+++

box-ball system

Korteweg-De Vries

nonlinear Schrödinger

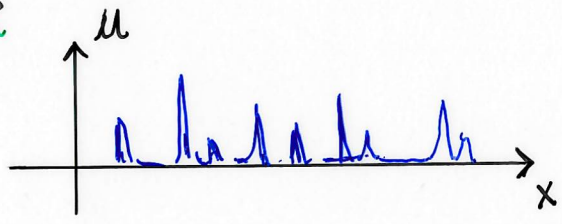
→ generalized Gibbs ensembles ← static

→ local GGEs \rightsquigarrow generalized hydrodynamics dynamics

// GGE averaged currents //

Euler hydrodynamics

- soliton gas



soliton-based G. El (2003)

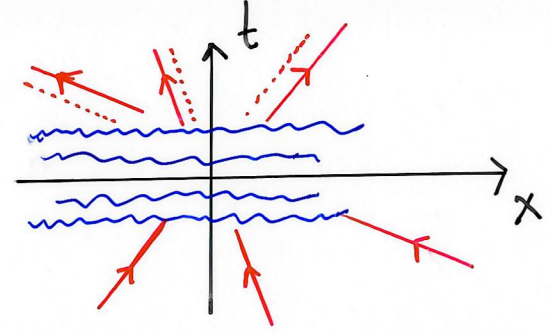
- particle system quantum / classical

// particles are solitons //

scattering shift

particle-based

Doyon et al, DeNardis et al (2016)



⇒ Toda lattice

$$H = \sum_j \left(\frac{1}{2} p_j^2 + e^{-(q_{j+1} - q_j)} \right)$$

two-particle
two-soliton

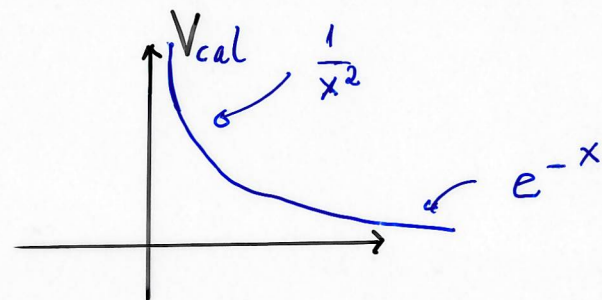
} scattering



2 distinct hydrodynamic scales

1. Calogero fluid *classical* $q_j, p_j, j=1, \dots, N$ on \mathbb{R} (1975)

$$H_N = \sum_{j=1}^N \frac{1}{2} p_j^2 + \sum_{2 \leq j=1}^N V_{cal}(q_i - q_j)$$

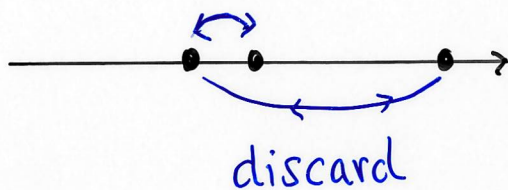


$$V_{cal}(x) = \frac{1}{\sinh^2(x)}$$

• no crossing $T_N = W_N \times \mathbb{R}^N$ Weyl $W_N = \{q_1 < \dots < q_N\}$

→ high density $V_{cal}(x) = \frac{1}{x^2}$ rational CMS (ideal fluid)

→ low density



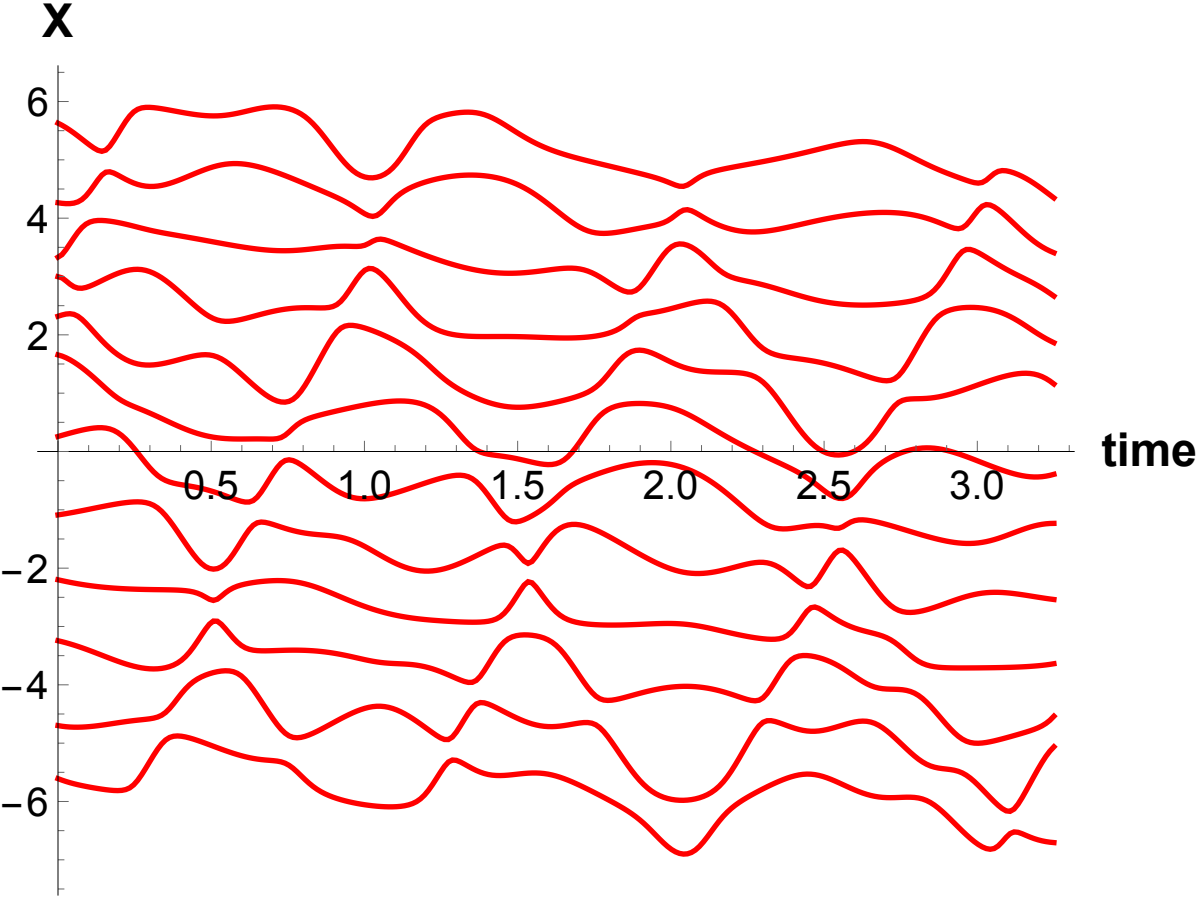
$$\Rightarrow H_{\text{ Toda, } N} = \sum_{j=1}^N \frac{1}{2} p_j^2 + \sum_{j=1}^{N-1} e^{-(q_{j+1} - q_j)}$$

stretch > 0

(free volume)

distinct < 0

Toda (1967)



2. Lax matrix

$N \times N$

$$L_{ij} = \delta_{ij} p_j + i (1 - \delta_{ij}) \frac{1}{\sinh(q_i - q_j)} \quad L = L^*$$

Lax pair $L(q, p), M(q, p) \rightsquigarrow \frac{d}{dt} L = [L, M]$

$\rightsquigarrow L \psi_\alpha = \lambda_\alpha \psi_\alpha$ eigenvalues are conserved **NON LOCAL** $\rightsquigarrow \left(\sum_{j=1}^N p_j^2 \right)^2$

• local fields **density** \rightsquigarrow total

$$Q^{[n]}(x) = \sum_{j=1}^N \delta(x - q_j) (L^n)_{jj} \quad Q^{[n]} = \int dx Q^{[n]}(x) = \text{tr}[L^n]$$

particle number $Q^{[0]}$, momentum $Q^{[1]}$, energy $Q^{[2]}$, ...

2nd Lax matrix

$$\frac{1}{\sinh} \Rightarrow \text{coth}$$

|| uniqueness of Lax ||

construction of 1-, 2-solitons

Kulkarni et al (2017)

3. GGE, generalized free energy

|| wish list ||

$$V_{cal, \ell}(x) = \sum_{m \in \mathbb{Z}} V_{cal}(x + m\ell)$$

double periodic Weierstrass



ℓ ring

Lax pair ✓ || more tricky ||

$$Q_{\ell}^{[n]} = \text{tr}[(L_{\ell})^n]$$

↪ GGE total charges $Q_{\ell}^{[n]}$

$$e^{-\sum_{n=0}^{\infty} \mu_n Q_{\ell}^{[n]}} = e^{-\text{tr}[V(L_{\ell})]}$$

relative to $dq^N dp^N$ on T_N

GGE parameters: $\frac{\ell}{N} = \nu, V$

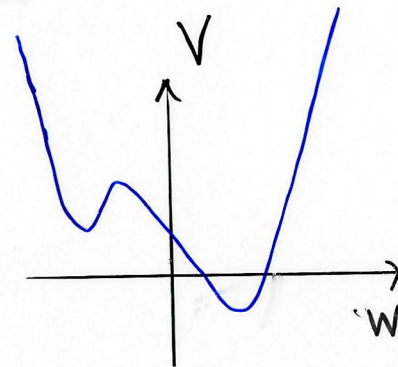
confining potential: $V(w) = \sum_{n=0}^{\infty} \mu_n w^n$ ||
 → momentum space ←

- free energy

$$\lim_{\ell \rightarrow \infty} -\frac{1}{\nu} \frac{1}{N} \log Z_N(\nu, V) = F(\nu, V)$$

|| functional of V ||

ideal gas
 $\prod_{j=1}^N e^{-V(p_j)}$



4. Density of states

$L_{N,l}$ random matrix under GGE, DOS

$$\rho_{Q,N}(\omega) = \frac{1}{N} \sum_{j=1}^N \delta(\omega - \lambda_j)$$

eigenvalue

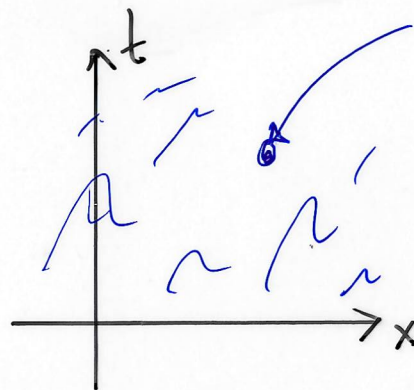
→ law of large numbers

$$\lim_{N \rightarrow \infty} \rho_{Q,N} = \rho_Q \leftarrow \text{deterministic}$$

$$\int_{\mathbb{R}} d\omega \rho_Q(\omega) \omega^n = \langle Q^{[n]}(0) \rangle_{GGE}$$

hydrodynamic picture, Lax filter

* is governed by Euler equations



$\rho_Q(x, t; \omega)$ *
local DOS, ν

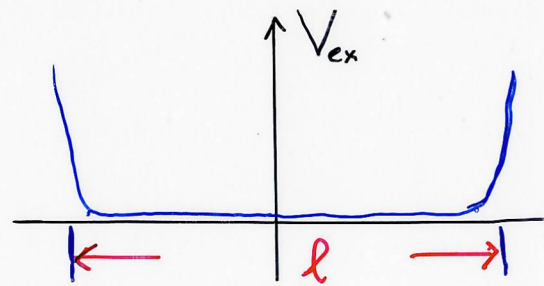
|| L filters slow degrees of freedom ||

5. External potential

$$T_N = W_N \times \mathbb{R}^N$$

↑

$$\{x_1 < \dots < x_N\}$$



$$Z_N(\nu, V) = \int_{T_N} d^N q d^N p e^{-\text{tr} V(L)} e^{-\sum_{j=1}^N V_{ex}(q_j)}$$

Ruijsenaars (1988 - 1995)

scattering coordinates

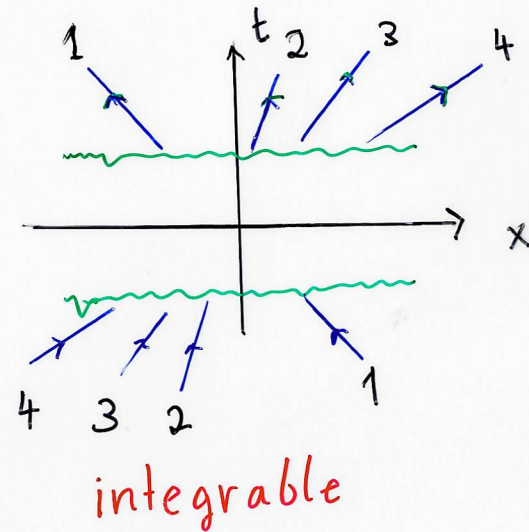
$\lambda \in W_N$ eigenvalues of L

$$\lim_{t \rightarrow \infty} p_j(t) = \lambda_j, \quad \lim_{t \rightarrow \infty} q_j(t) - \lambda_j t = \phi_j \in \mathbb{R}$$

$\Phi : (\lambda, \phi) \mapsto (q, p)$ one-to-one, canonical
generic

Result (R 1988) algebraic construction of Φ

$$Z_N(\nu, V) = \int_{T_N} d^N \lambda d^N \phi e^{-\sum_{i=1}^N V(\lambda_i)} e^{-\sum_{i=1}^N V_{ex}(q_i(\lambda, \phi))}$$



• special choice $V_{ex}(x) = e^{-\ell/2} \cosh x$

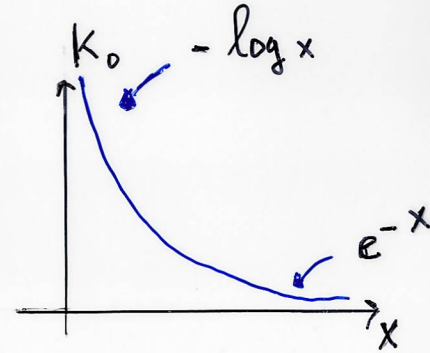
$$\sum_{j=1}^N V_{ex}(a_j) = \sum_{j=1}^N e^{-\ell/2} Y_j \cosh \phi_j$$

$$Y_j = \prod_{\substack{m=1 \\ m \neq j}}^N \left(1 + \frac{1}{(\lambda_m - \lambda_j)^2} \right)^{1/2}$$

confining!

modified Bessel

$$\Rightarrow Z_N(z, V) = \frac{1}{N!} \int_{\mathbb{R}^N} d\lambda \prod_{j=1}^N e^{-V(\lambda_j)} \prod_{j=1}^N 2K_0(2e^{-\ell/2} Y_j)$$



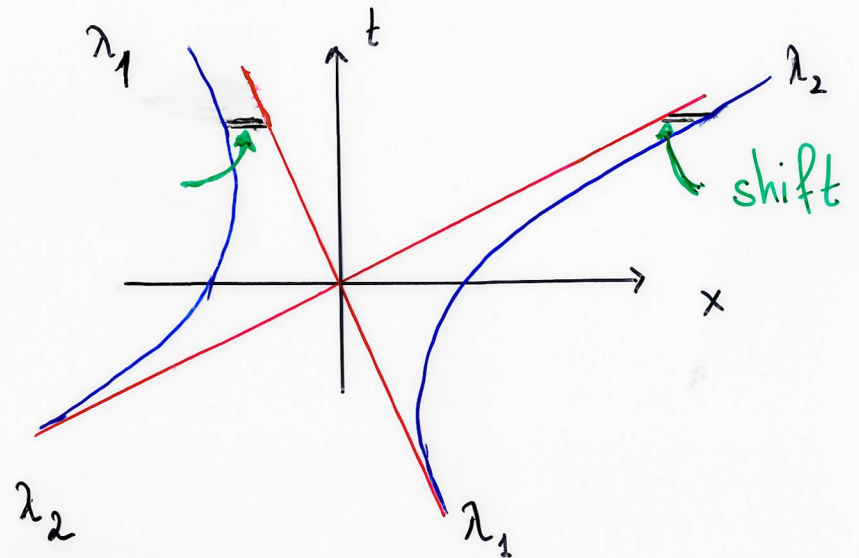
|| mean-field ||

• Calogero 2-particle scattering shift

$$\phi_{ca}(w) = -\log\left(1 + \frac{i}{w^2}\right)$$

$$w = \lambda_1 - \lambda_2$$

|| quasi-particle ||



6. free energy functional (1-particle)

$$\rho \geq 0, \int_{\mathbb{R}} dw \rho(w) = \frac{1}{v}$$

$$F(\rho) = \int_{\mathbb{R}} dw \rho(w) \left(V(w) - 1 + \log \rho(w) - \log \left(1 + \int_{\mathbb{R}} dw' \rho(w') \phi_{ca}(w-w') \right) \right)$$

minimizer ρ^* (unique)

$$F(\rho^*) = F(v, V) = \lim_{N \rightarrow \infty} -\frac{1}{N} \log Z_N(v, V) \quad \parallel \quad v = \frac{l}{N}$$

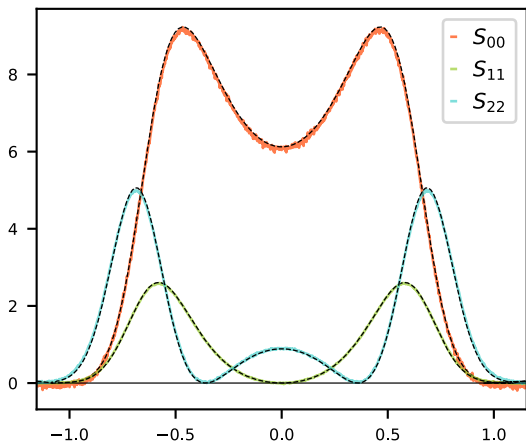
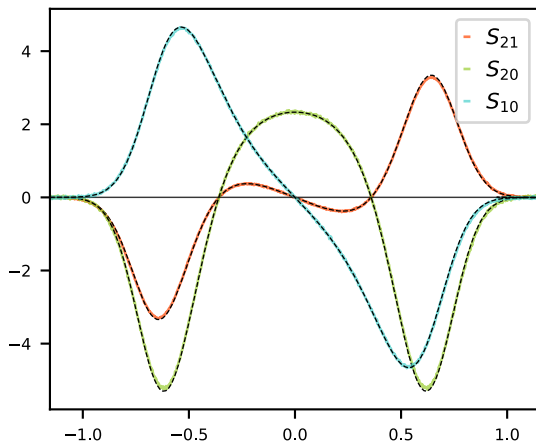
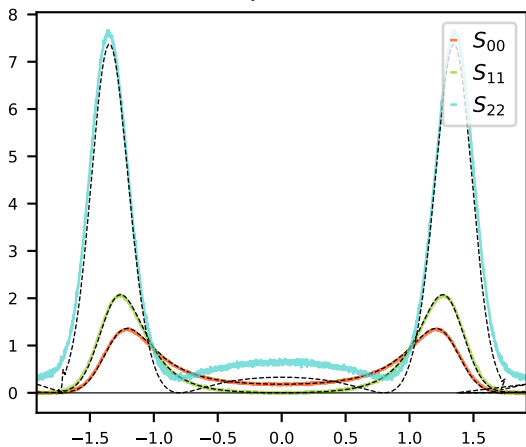
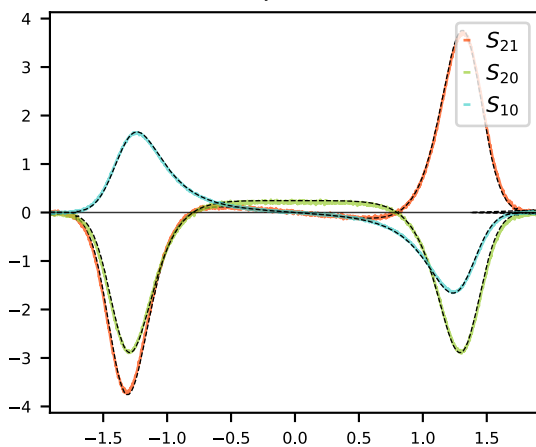
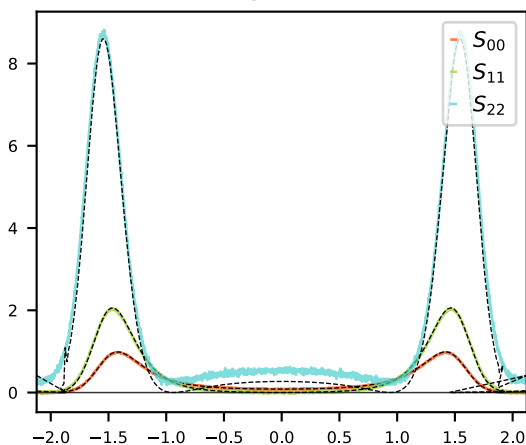
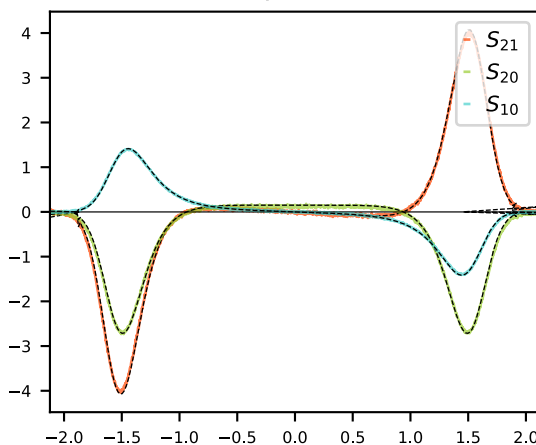
$$\text{Lax DOS} \quad \rho_a = v \rho^*$$

hydrodynamic equation:

GHD

$$\partial_t \rho(x, t; w) + \partial_x \left(v^{\text{eff}}(x, t; w) \rho(x, t; w) \right) = 0$$

$$v^{\text{eff}}(w) = w + \int dw' \rho(w') \phi_{ca}(w'-w) \left(v^{\text{eff}}(w') - v^{\text{eff}}(w) \right)$$

$S_{00}, S_{11}, S_{22}, \beta = 0.5, P = 0.32$  $S_{21}, S_{20}, S_{10}, \beta = 0.5, P = 0.32$  $S_{00}, S_{11}, S_{22}, \beta = 0.5, P = 0.95$  $S_{21}, S_{20}, S_{10}, \beta = 0.5, P = 0.95$  $S_{00}, S_{11}, S_{22}, \beta = 0.5, P = 1.21$  $S_{21}, S_{20}, S_{10}, \beta = 0.5, P = 1.21$ 

Outlook

- GGE spatial mixing for Toda Mazzuca, Memin (2023)
exponential
- MD density of states ✓
- rational $\frac{1}{x^2}$ scattering shift = 0

$$\Rightarrow v^{-ik}(w) = w$$

$$\partial_t \rho + \partial_x v \rho = 0$$